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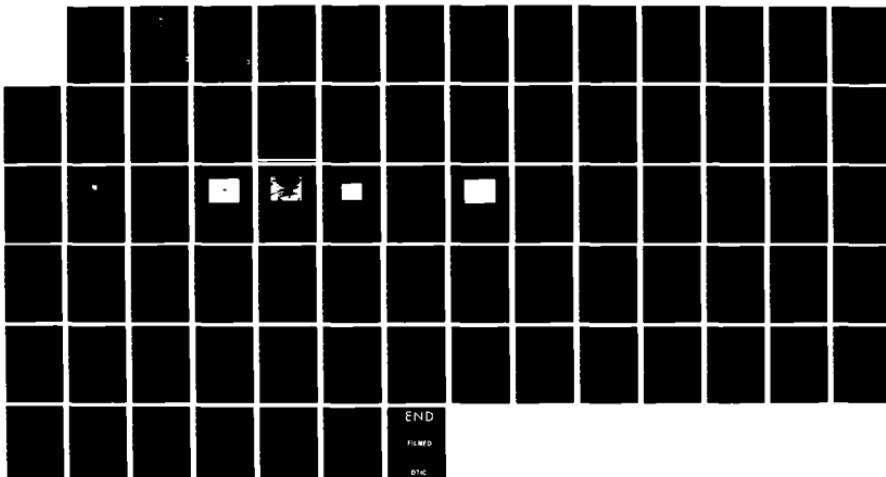
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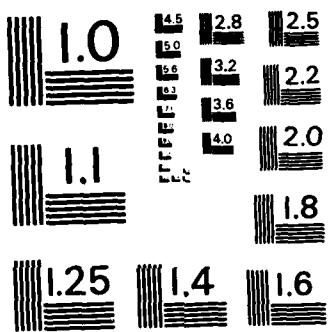
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Dall-Null Tester for
Spaceborne Applications

Thesis

Randy L. Wingler
Captain, USAF

AFIT/GEP/PH/84D-14

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Dall-Null Tester for Spaceborne Applications

Thesis

Presented to the Faculty of the School of Engineering
of the Air Force Institute of Technology

Air University

In Partial Fulfillment of the
Requirements for the Degree of
Master of Science in Engineering Physics

Ronald L. Winkler, B.S.

Capitol City, D.C.

July 1984

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Preface

The purpose of this study was to design an adaptive optic system for a space based telescope. The design is built around a Dall-Hull tester. Much of the fine detail of the system was not researched because of the great difficulty and time required to correct all of the previous theory on which the design is based. However, the basis of a complete system is covered.

The making of the Focogram in the experimental section was very enjoyable and a nice break from the pitfalls of straight theory work. The results from the experiment showed the errors of the previous theories and paved the way to correcting them.

Several people have assisted me in this project. They are my advisor Dr. Doug Shantz, who always kept me headed in the right direction, but made no lock and question every bend along the way; Paul McCarty, the librarian who tracked down all the references that I requested; my class mate, Captain David C. Ladd, who wrote a lot of computer programs and helped me to get it. Gertts J. Hiltz, who helped me gather and organize all the data for the final design. Paul M. Johnson, who helped me to make the drawings for the Dall-Hull tester. I would like to thank all the people at the University of Alberta, especially the faculty, for their help and support.

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Abstract

This is a study to design a self-correcting primary mirror system for a space telescope. The design is centered around a Dall-Null tester (a Foucault knife-edge tester with compensating lens). An in-depth study of the theory of the Foucault test from Foucault's original publications to current work is presented. Also shortcomings of the diffraction approach are shown. The findings of a simple experiment showed the way to the correct explanation as to the workings of the test. Based on this new explanation, a computer program to find the error in the surface of the mirror from the irradiance pattern provided by the Dall-Null tester was developed. The computer program with a sample run is included in the appendices A and B.

The basic design of an adaptive optic system for a space-borne application is also presented in the paper. This design has the desired quality of being able to correct the mirror while the telescope is in use. The equations being independent of wavelength allow the heating to be applied to any wavelength or at any time. This will reduce the time required to correct the mirror.

Chapter 1

Introduction and Background

Ever since Galileo turned his telescope to the heavens, man has tried to improve this sky gazing instrument in order to gain more information from the skies. The best improvement that can be made in any telescope is to increase its light gathering capabilities. Increasing the size and improving the quality of the main optical element are the easiest ways to improve the light gathering capability of a telescope. The primary, or main optical element in a reflecting telescope, is the main mirror. And in a reflecting telescope it is the large fixed lens. Since most large telescopes are reflectors, due mainly to their compact size and the lighter weight of the system, this paper will concern itself with improvements that are more suitable for reflecting telescopes.

The main optical element of the primary mirror in a telescope is the only way of increasing the light gathering capability. However, it is important when, for it may be either a large telescope or a small one, to increase the size of the primary mirror. This is because the amount of light gathered by the telescope is proportional to the area of the primary mirror. The area of a circle is proportional to the square of the radius. Therefore, if the radius of the primary mirror is doubled, the area is increased four times. If the radius is tripled, the area is increased nine times. This is why the larger reflecting telescopes have larger primary mirrors. The larger the primary mirror, the more light is gathered by the telescope.

It is the reflective coating on the surface of the mirror that determines the reflectivity of the mirror. Therefore, any improvement in the reflectivity of the mirror is limited to the thickness and type of coating applied. Since the reflectivity is set in this manner, the only way left to improve the quality is to rid the surface of the mirror of defects. Most surface defects or imperfections are removed while the mirror is being fabricated.

Optical telescope mirrors are made by grinding glass into the desired shape and then coating them with a highly reflecting substance, like silver. As the glass is ground into the desired shape it is visually monitored for defects and imperfections. The most commonly used test for checking the surface is the Foucault knife-edge test (or one of its derivatives). The Foucault test has one large difference. It is designed to check only the flat surfaces and not the spherical surfaces that are created during the fabrication. In the Foucault test, a thin, flat, rectangular, parallel edge, such as a pocket knife blade, is held at a constant angle to the surface of the mirror. The light reflected off the surface is focused onto a screen. If the surface is flat, the image will be sharp and clear. If the surface is not flat, the image will be blurred and distorted.

The Foucault test is not the only method used to check the surface of a mirror. Another method is the interferometer. An interferometer uses a laser beam to create interference patterns. These patterns are then analyzed to determine the quality of the mirror's surface. Interferometers are more accurate than Foucault tests because they can measure much smaller surface irregularities. They are also faster and easier to use. However, they are more expensive and require more specialized equipment.

the results of the usual test of a perfect spherical surface. This method of testing an optical surface using the proper lens in front of the Foucault tester now is known as the Dall-Null test. If the focal length of the mirror, the focal length of the desired lens, and the desired focal ratio of the mirror are known, then the proper compensating lens can be determined from a nomograph like the one found in *Sky and Telescope* (6:212).

Since the Dall-Null test is nothing more than a Foucault test with a lens in front of the knife-edge, all the theory that explains the Foucault test also explains the Dall-Null test. The theory of the Foucault test as it developed from Foucault's original publication, to the present, is presented in Chapter 6. To implement the use of the theory as presented in that chapter, one must be able to detect an incorrect deflection in the profile of the mirror. This can be done as follows.

In order to implement the theory of the Foucault test, one must have the proper equipment. The first item of equipment required is a Foucault tester. This is a device which projects a beam of light through a lens onto a mirror. The beam is reflected back through the lens and focused onto a screen. The beam is focused at a point on the screen which is directly opposite the point where it was projected. This is called a real image.

The second item of equipment required is a lens. This lens must be a converging lens. It must be a lens which has a focal length greater than the distance between the lens and the mirror. This is called a diverging lens.

The third item of equipment required is a mirror. This mirror must be a plane mirror. It must be a mirror which reflects light. This is called a reflecting mirror.

The fourth item of equipment required is a knife-edge. This knife-edge must be a sharp edge. It must be a edge which cuts through the light. This is called a cutting edge.

The fifth item of equipment required is a lens. This lens must be a diverging lens. It must be a lens which has a focal length less than the distance between the lens and the mirror. This is called a converging lens.

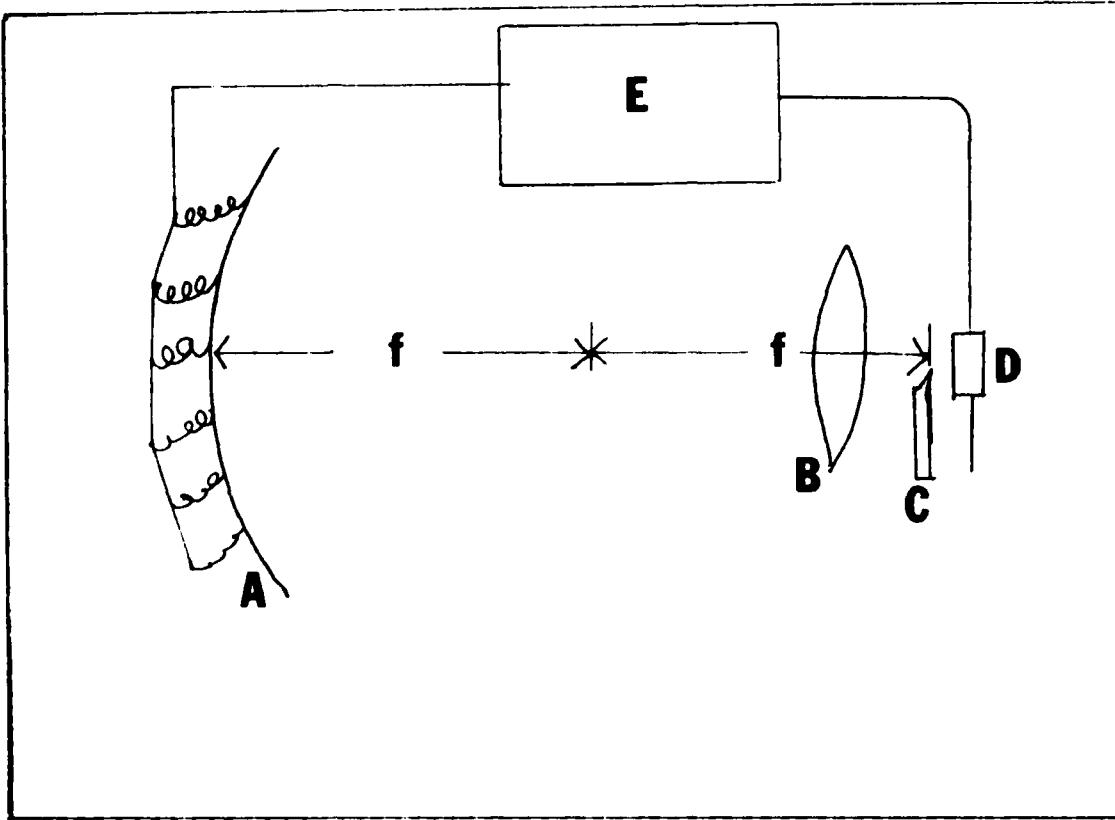


Figure 1 The basic layout of a Dall-Null tester
 A is the deformable mirror with attached actuators
 B is the compensating lens
 C is the knife-edge
 D is the detector for reading the irradiance
 E is the computer and other electronic equipment

to escape the atmospheric limitations. This does cause some problems. One such problem is, "How can the surface of the mirror be tested to see if it has been twisted and warped out of shape by thermal and gravitational gradients as it spins around the earth?". This leads to the second question, "If surface defects do occur, how are they detected and corrected?".

One way to answer both of the above questions is to make the telescope with an adaptive optics system. This adaptive optics system consists of the following: a deformable main mirror with actuators, a computer, some electronics, a Dall compensating lens, and two Foucault testers. Figure 1 shows the basic layout of an adaptive optics system of a telescope. The deformable mirror is made by shaping a semiflexible material into the desired shape, and then coating it with a highly reflective substance (17,28). Then actuators (devices that can apply a push or pull force) are mounted on the back surface of the mirror. There are two major types of actuators. The first is the force type that uses either a stepping motor, or an electrodynamic voice coil to generate a push/pull force. The second is the piezoelectric type that contracts or expands with applied voltage (1). The computer's job is to calculate the error of the surface of the mirror from the irradiance of the images caused by each of the Foucault testers, and then to take the error information and correct the surface of the mirror through movement of the actuators. Two Foucault testers are needed to produce both a vertical and a horizontal scan. The two knife-edges are located approximately at the center of curvature, but orientated at 90 degrees from each other. The proper Dall compensating lens corrects the Foucault test for parabolic surfaces (6,7). The additional electronic equipment needed to complete a system consists of a sensor to read the irradiance, equipment to digitize the irradiance

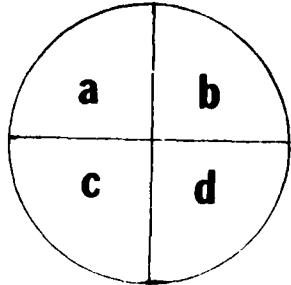


Figure 5. Mirror divided into four segments

pattern that is small ($\approx \lambda$). (Figure 4). These Airy patterns would either be passed by the knife-edge, or stopped completely, with just a few patterns left to straddle the knife-edge. It is these straddling patterns that would produce the effect of equation (2). But, because the light is incoherent, the irradiance of the waves add and not just their fields. This removes the singularities.

Now if the mirror was segmented into sections, as in Figure 5, each segment having a different local slope, the position of the light at the focus would be determined by the slope of each segment. Figure 6 shows the relative position of the light from each segment of the mirror relative to the knife-edge. If the slope is that of a perfect surface the knife-edge would split the image, case b, passing half and stopping half. If the slope was such that it raised the

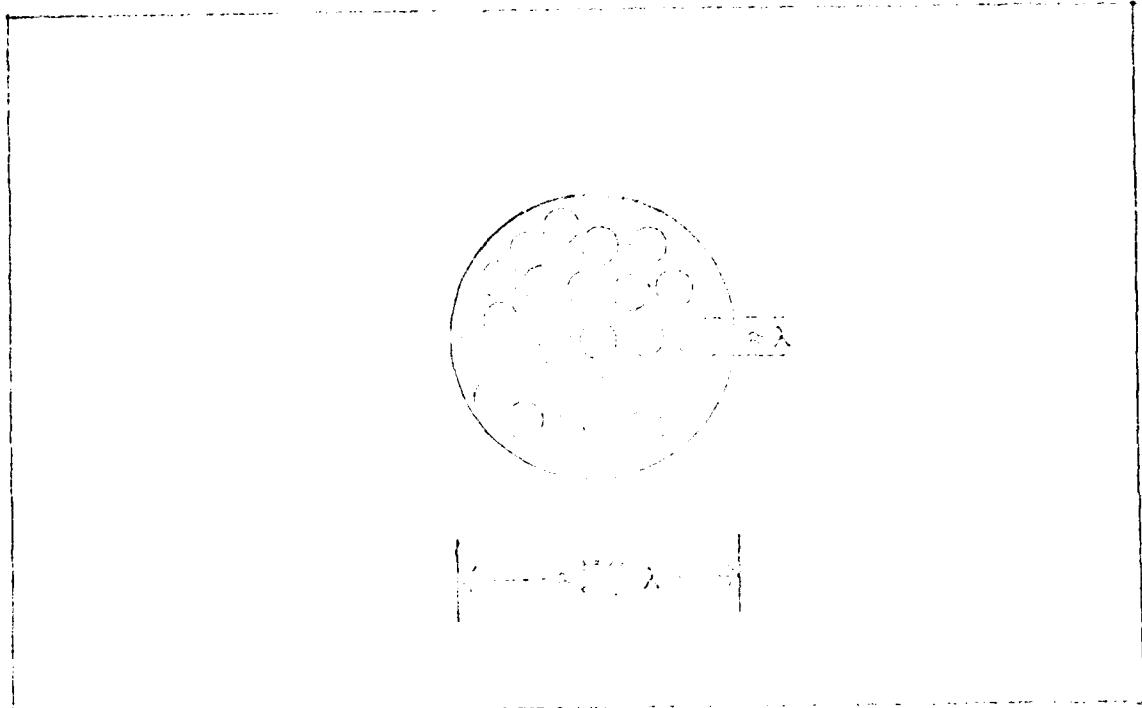


Figure 4. The front propagation at the time of the front at the local period

$A(\underline{x})$ = the aperture function

$$\underline{x} = (x_1, x_2)$$

$$\underline{y} = (y_1, y_2)$$

$$\underline{z} = (z_1, z_2)$$

b = the distance the knife is displaced from the optic axis

But when equation 2 is used (being careful in expanding out all transforms and exponentials) to relate the irradiance of the image to the errors of the mirror, the resulting equations are so messy and complex they can not be solved. This is one of reasons the experiment described in chapter three was conducted.

The experiment described in Chapter Three was performed to verify what kind of results are obtained from an actual Foucault test, and to examine some of the operating parameters. The results of the experiment showed no existence of a bright rim around the image, and that the aperture did not affect the image at the screen as long as it does not intersect the beam. The experiment also showed that the light source had to be incoherent, and that monochromatic light was better than plain white for testing lenses due to the different focal length of each wavelength.

When using the laser as the source of radiation in the experiment, it became apparent that the diffraction theory of the Foucault test was inadequate to describe the results.

Where a = slit width

$$= M\lambda R/b$$

$$n^o = m\lambda R/b$$

b = pupil radius

M = number of Airy patterns across slit

m = number of Airy patterns the lower edge is displaced from the center of the image

R = the radius out to the image from the knife-edge

λ = wavelength

The only real difference between this equation and that of Linfoot's is that the factor $1/(y-y')$ which gave the singularity is now replaced by the sinc function which has no singularity.

Welford's equation is basically the same result that the author, working in conjunction with Dr. Blankland, obtained upon applying Huygen's principle to the problem, (Figure 3). The resulting equation for the resultant amplitude, $P(z)$, is

$$P(z) = B \int_{-\infty}^{\infty} K(x, z) f(x) dx \quad (2)$$

where

$$K(x, z) = \int_{-\infty}^{\infty} dz' e^{-jz'x} e^{-jz'k^2/2} \left[1 - \frac{f(z')}{f(z)} \right] \left(\frac{1}{z'} - \frac{1}{z} + \frac{1}{2} \frac{d}{dz'} \left(\frac{1}{z'} - \frac{1}{z} \right) \right)^{-1/2}$$

$$B = C_0 \pi^3 \lambda^3$$

$$C_0 = 0.476 \times 10^{-12} \text{ rad}^2 \text{ nm}^2 \text{ cm}^2 \text{ sec}^{-1}$$

$$C_0 = 1.05 \times 10^{-12} \text{ rad}^2 \text{ nm}^2 \text{ cm}^2 \text{ sec}^{-1}$$

$$\lambda = 0.555 \text{ microns} \quad k = 2\pi/\lambda = 11,300 \text{ cm}^{-1}$$

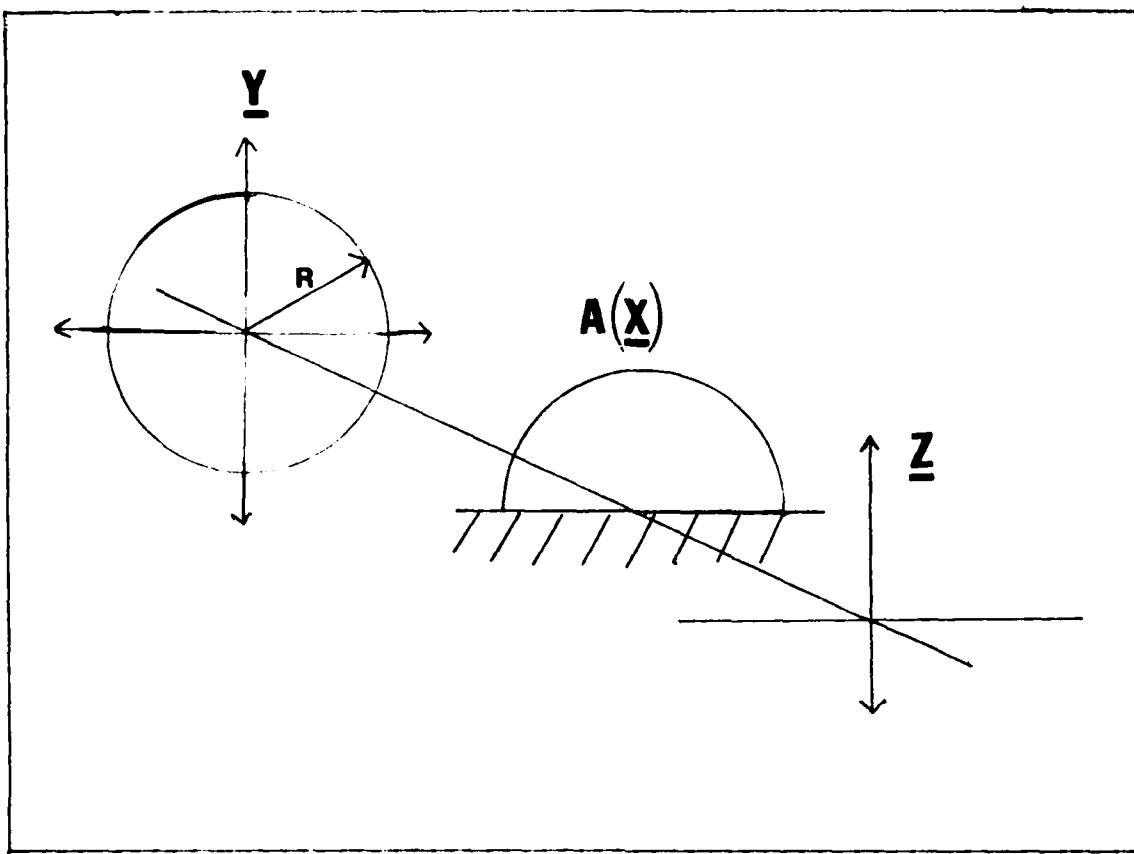


Figure 3 Orientation for the application of Huygens' Principle

disregard it as there could not be an actual infinity in the complex amplitude. Linfoot said that this infinity was due to the assumption of an infinite aperture. Welford decided to test this, and rederived Linfoot's work with a proper aperture function included. He used a very wide slit in the place of the knife-edge. Using Figure 2 as a reference, the equation Welford derived for the complex amplitude is:

$$U(x', y') = \int_{-\infty}^{\infty} a \operatorname{sinc}(M/b(y-y')) \exp(-\pi i/b(M+2m)(y-y')) F(x', y) dy$$

A plot of the perfect mirror's irradiance is found on page 140 of reference 23, and the plot of the perfect cassegrain mirror is on page 141. The point to note here is that the graphs and theory predict irradiance well out into the surrounding area outside of the mirror's boundary. This is not observed in practice.

W. T. Welford wrote an article (33) on the theory of the Foucault test in 1970, in which he noted that in Linfoot's, and his predecessors', computations a bright rim was predicted around the edge of the pupil of the system under test. Users of the test do not observe this bright ring, and

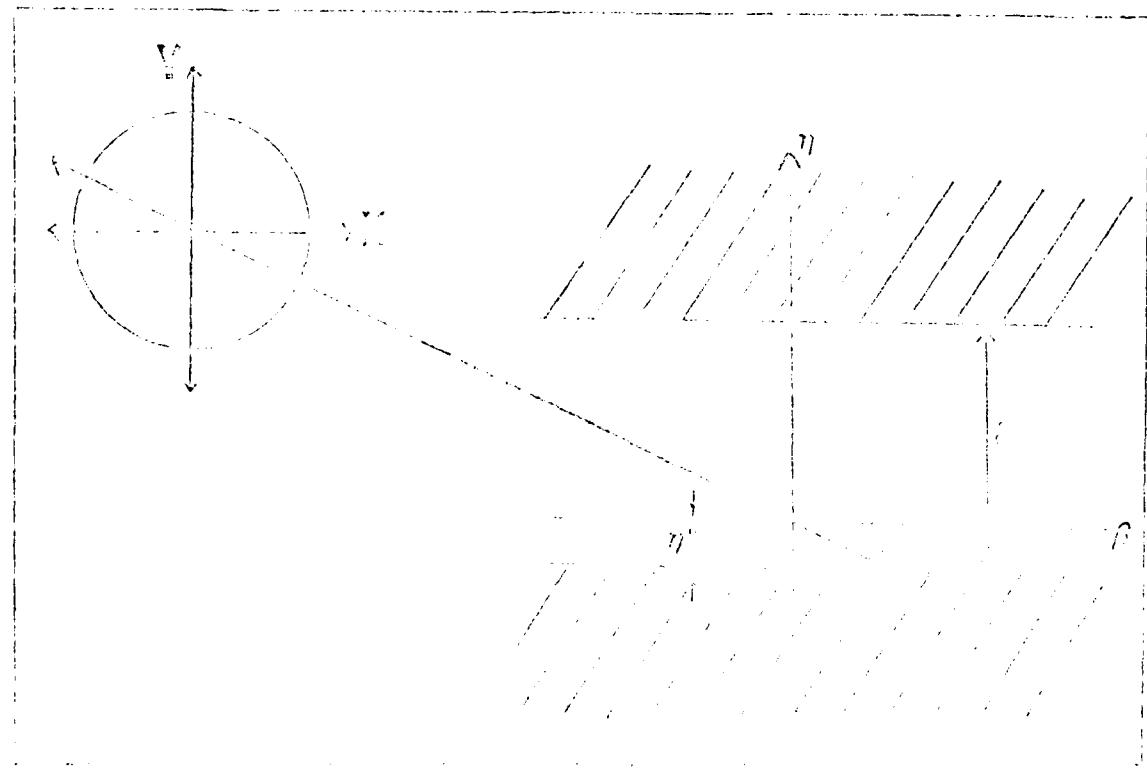


Figure 1.1.1. A schematic diagram of the Foucault test.

using a one inch diameter lens to expand the image; the "sec" term is less than one part in 100,000 at the edge of τ , where τ in this case is the semicircular edge corresponding to the

value of $[u^2 + v^2]^{1/2}$ which is approximately 1300 (inch)^{-1} . Disregarding the effects of the finite aperture of the viewing system, one can replace τ with infinity as the image is defined to be zero outside of the mirror. The complex amplitude now can be written,

$$D(x', y') = \frac{1}{2\pi} \int_0^\infty du \int_{-\infty}^\infty \exp(-i\omega u^2 - i\omega v^2) W(u, v) dv$$

This reduces to the expression:

$$2\pi D(x', y') = n E(x', y') + i \int_{-\infty}^{\infty} f(x', y') e^{i\omega v^2} dv$$

and the observed irradiance is:

$$I(x', y') = n^2 \left| D(x', y') \right|^2$$

The constant n is given by the equation:

$$\{f(x', y')\} = \frac{2}{\pi} \int_0^\infty \left[\int_{-\infty}^\infty \left| W(u, v) \right|^2 dv \right] du$$

This is saying the irradiance is uniform across the surface and zero outside the boundary of the mirror. If the surface of the actual mirror surface is labeled M , which is within a few wavelengths of a true spherical surface labeled M^0 , then the wave at any point (x, y) is:

$$E(x, y) = |E(x, y)| \exp((-2\pi i/\lambda) \theta(x, y))$$

where $|E(x, y)|$ = the amplitude at the point (x, y)

$\theta(x, y)$ = the phase at the point (x, y)

λ = the wavelength

Now applying Huygen's principle along with the Paraxial approximation twice, once to the half-edge and again to the image, results in the equation for the complex amplitude (the square of the amplitude is the irradiance) as:

$$D(x', y') = \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \exp(-i\pi(u' - v')) U(u, v) \exp\left(\frac{-2\pi i}{\lambda}(u' + v')\right) du dv$$

$$\text{where } u = 2\pi x / \lambda$$

$$v = 2\pi y / \lambda$$

π = The constant π = 3.141592653589793 to 15 decimal places

$$U(u, v) = \frac{1}{\pi^2} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} E(x, y) \exp(-i\pi(u + v)) dx dy$$

$$E(x, y) = |E(x, y)| \exp((-2\pi i/\lambda) \theta(x, y))$$

$$\theta(x, y) = \theta_0 + k(x - x_0) + k(y - y_0)$$

$$\theta_0 = \text{constant}$$

Where R is some retardation function

x = the position out from the center

(the radius has been normalized to one and y is just a variable of integration)

In the case of the perfect mirror, the retardation function, R , is equal to zero everywhere, and the equation reduces to Rayleigh's result, infinity at the edge of the mirror.

S. C. B. Gascoigne, working with R. H. Linfoot, took Zernike's work as his starting point and derived expressions for the irradiance that are valid for errors of arbitrary form and amount, in terms of definite integrals (Hilbert transform). Gascoigne's equation for the irradiance of a perfect mirror is equivalent to Zernike's (equation 1). Gascoigne has tabulated his irradiance results on page 332 of reference 12 and plotted the results on page 334 of same reference.

It wasn't until Linfoot published his work in 1955, before anyone worked out the theory completely for two or more dimensions (33). The equations will be much more uniform if the solution is in the form of the radial coordinate:

$$I(r, \theta, \phi) = 1 + \sum_{n=1}^{\infty} [C_n(r) J_n(\theta) Y_n(\phi)]$$

$$J_n(r) = \int_0^\infty r^{n+1} J_n(r') e^{-r' r} dr'$$

was well out of the focal plane. His calculated irradiance ratios were then matched against photographic observed ratios with less than ten percent difference between corresponding points. Thus, Rayleigh's theory stood its ground, untouched until 1934.

In 1934, professor F. Zernike published his diffraction theory of the Foucault test(35). Zernike felt that the geometrical optic approach to the Foucault test did explain things well enough for the lens grinder to find the spots where more grinding was necessary, but in a test where wavelength defects could be detected, only wave optics and diffraction could completely describe the actual happenings and give the test's ultimate limit of sensitivity (35:377). He envisioned the surface of the mirror as a diffraction grating that would give first order spectra on both sides of the central image, which he called a phase grating. He built up his work by expressing everything in terms of this phase grating. From his phase grating expressions, he developed what is now called the method of phase contrast. This phase contrast method assigns phases to the different shades of irradiance in the image. Zernike's equation for the observed intensity, $I(x)$, is:

$$I(x) = \pi^2 + \ln^2 \left| \frac{1-x}{1+x} \right| - 2\pi \int_{-1}^1 \frac{R(y) - R(x)}{y-x} dy \quad (1)$$

$$I(x) = \pi + \ln^2 \left| \frac{D - x}{D + x} \right|$$

for the irradiance inside the edge of the mirror and

$$I(x) = \ln^2 \left| \frac{D - x}{D + x} \right|$$

for the irradiance outside the edge.

where D = the radius of the mirror

x = the position out from the center

At the edge of the mirror, when $x = D$, the logarithmic infinity that results is what Rayleigh said accounted for the bright ring around the image.

Sudhansukumar Banerji, in 1918, wrote that he had observed irradiance fluctuations of the entire surface when the knife-edge was considerably advanced and that there were no known repetitions or continuation of this effect (2). He then proceeded to show that the fluctuations could be explained using Rayleigh's theory. Banerji then pointed out that Rayleigh had difficulties finding the case where the knife-edge was exactly at the focus, and so proposed the possibility of the fluctuations occurring around the "focal" and "antifocal" points. He also found that the fluctuations would increase as the distance between the source and the screen increased. He also found that the fluctuations were greater for smaller apertures and for larger distances between the source and the screen. He also found that the fluctuations were greater for smaller apertures and for larger distances between the source and the screen.

bright ring about the edge. The model also accounted for the bright line which appears at a step-discontinuity on the mirror surface.(27)

The theory of the Foucault test, as developed by Lord Rayleigh, gives irradiance of the field for the general case of any mirror, as viewed in the direction Θ (relative to the optic axis) in terms of the Fresnel Cosine (Ci) and Sine (Si) integrals as:

$$I = \left[Si\left[\frac{2\pi}{\lambda} \beta(1+\theta/\beta) \zeta_2\right] - Si\left[\frac{2\pi}{\lambda} \beta(1+\theta/\beta) \zeta_1\right] + Si\left[\frac{2\pi}{\lambda} \beta(1-\theta/\beta) \zeta_2\right] \right. \\ \left. - Si\left[\frac{2\pi}{\lambda} \beta(1-\theta/\beta) \zeta_1\right]\right]^2 + \left[Ci\left[\frac{2\pi}{\lambda} \beta(1-\theta/\beta)\right]_2 - Ci\left[\frac{2\pi}{\lambda} \beta(1-\theta/\beta)\right]_1 \right. \\ \left. - Ci\left[\frac{2\pi}{\lambda} \beta(1+\theta/\beta)\zeta_2\right] + Ci\left[\frac{2\pi}{\lambda} \beta(1+\theta/\beta)\zeta_1\right]\right]^2$$

where λ = Wavelength

ζ_2 = upper limit of aperture

ζ_1 = lower limit of aperture

θ = angular semi-aperture of the lens under test

Applying Rayleigh's irradiance equation for the special case of a perfect spherical mirror, tested at the center of curvature, with a change of variable in order to do the integration, the above equation for $I(x)$ reduces to:

Chapter 2

The Theory of the Foucault Test

After 1858 and 1859, when Leon Foucault published the first accounts of the Foucault knife-edge test, it rapidly became the optical test to verify the surface condition of a lens or a mirror (9,10,25:231). His first article described how to conduct the test, how to interpret the results, but not how or why it works (9). In his second article, he did attempt to explain the happenings in terms of geometric ray optics (10). This explanation worked well for extremely large errors but did not account for all of the observed effects. As Linfoot pointed out, the effects that are produced by the small errors (on the order of one to ten wavelengths) often caused trouble for the inexperienced mirror/lens grinder (24:128). One of these errors is the bright ring around the image of a properly ground lens. According to Foucault's original work, this bright ring could only be explained by saying it represented a steep narrow-turned edge, but there is no such edge on the surface of a properly ground lens.

In 1917, Lord Baron Rayleigh made the first attempt to derive a complete theory of how and why the Foucault test worked based on diffraction theory and wave optics. He used a simplified two-dimensional model (all equations and work were in one dimension) which qualitatively explained the

of the images for the computer, and equipment to provide the feedback between the actuators and the computer. Chapter four will discuss how the error in the surface height of the mirror is calculated from the irradiance patterns.

The complete system, deformable mirror, actuators, Dall-Hull tester, computer, and the computer software for reading the focogram and controlling the actuators, comprised the necessary equipment to make up a self-correcting mirror system for a space-borne telescope application. This system has one big advantage over other possible systems for telescope applications. It works at the center of curvature and not at the focus. This allows the testing and correction of the primary mirror to occur while the telescope is in use.

Chapter five will present conclusions of this study with recommendations for further study, the possibilities of adaptive optics applications, and experimental results, limitations, and possible future directions.

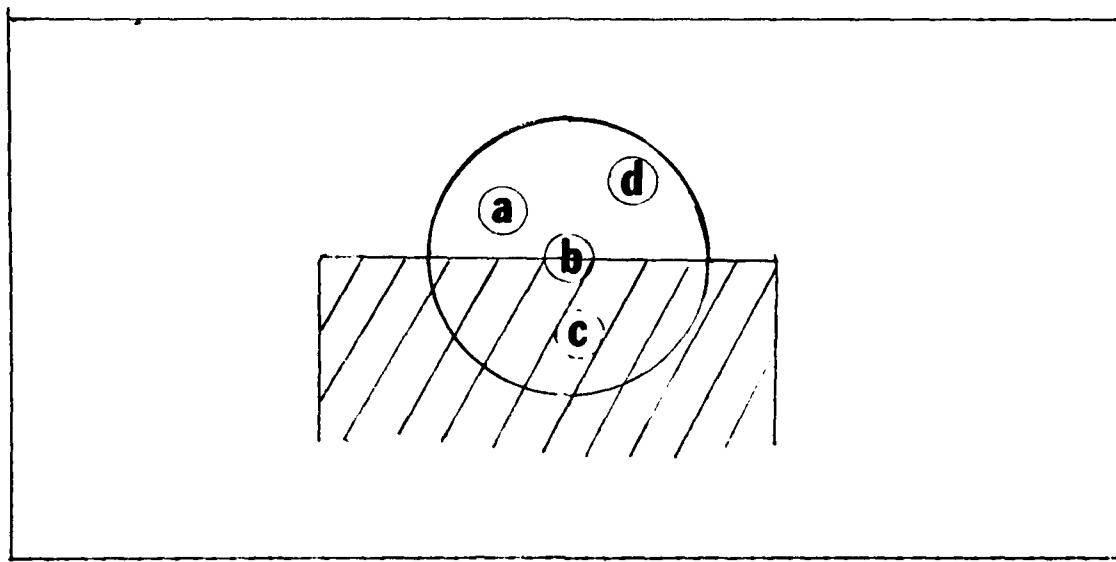


Figure 6. The relative position of the Arisy patterns at the knife-edge.

image, cases a and d, then all of the image will pass the knife-edge. But if the slope lowered the image, case c, the knife-edge would cut it off completely. Thus the irradiance of the resulting image would be determined by how much of each segment's image was passed by the knife-edge. Figure 7 shows the result for the current situation.

Thus the correct theory of the Foucault test is revealed, and likewise, so is the theory of the Dall-Null test. And in chapter four this theory will be put to work to find the error in the height of the mirror surface.

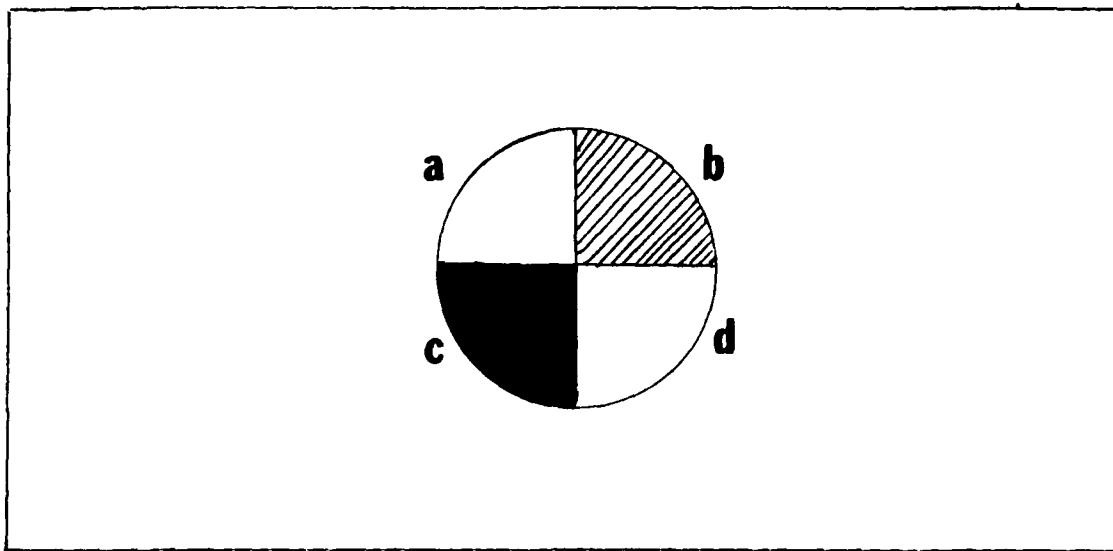


Figure 7. The image of the mirror

Chapter 3

An Experiment to Produce a Focogram

The Focograms presented in the literature (25:233-241, 20:390) did not show any bright rim about the edge that theory, according to Linfoot and his followers, says should be there. To answer the questions, "Is the theory right or is there something happening that is being missed?" and "Is Welford right in that it is the aperture function that will cure the infinity at the edge?", a simple experiment was set up to produce a Focogram of a lens.

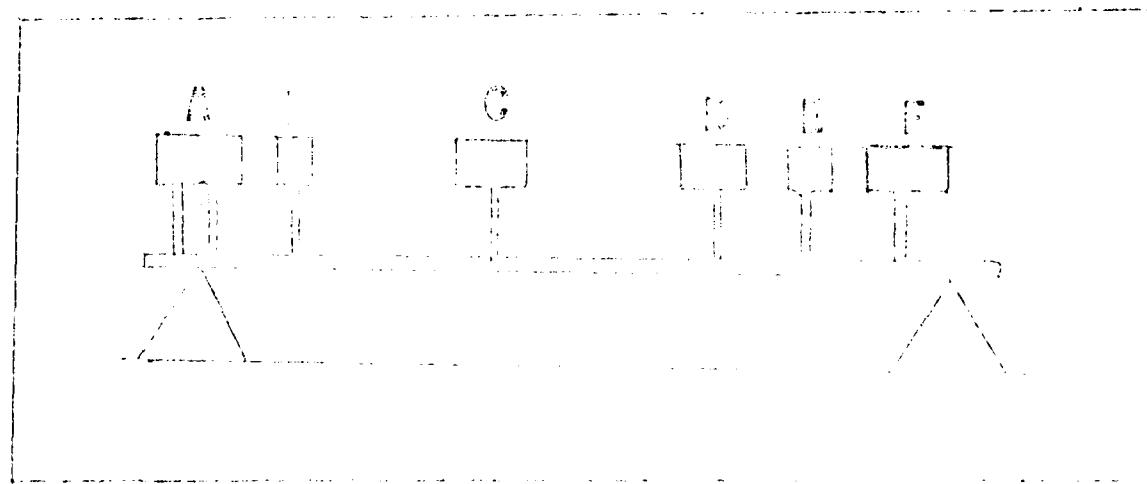


Figure 9. Experimental setup.

After the light from the source had passed through the lens, it was divided by the beam splitter. One part of the beam was directed to the camera obscura, and the other part continued to the second lens. This second lens was also imaged by the camera obscura. The two images were then compared.

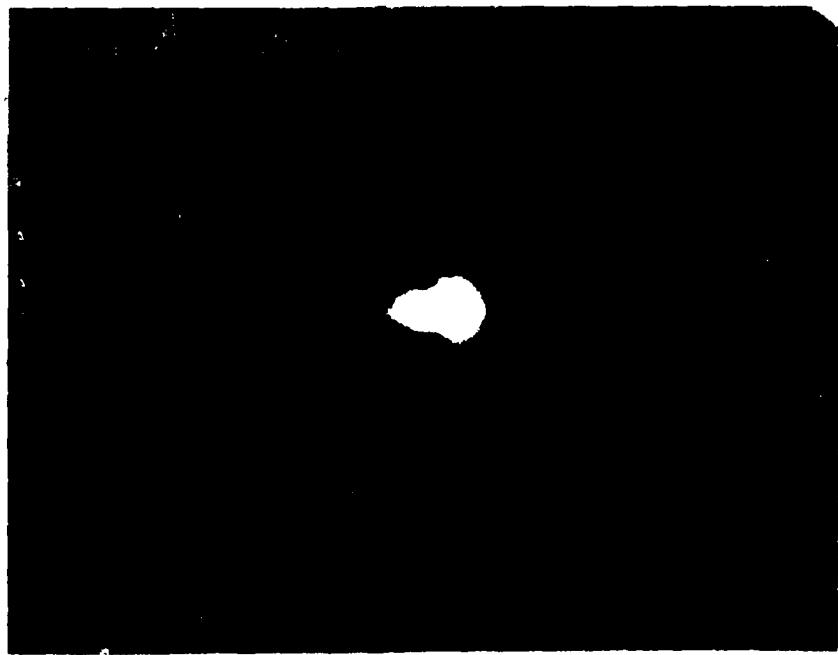


Figure 9. A typical example of a blank page from the original document.

Figure 9 shows a typical example of a blank page from the original document. The page contains several large black rectangular redaction boxes, which are used to obscure sensitive information. One such box is located in the center of the page, covering the majority of the content. Another smaller redaction box is located near the bottom left corner. The rest of the page is mostly blank white space, with some minor scanning artifacts visible along the edges.

The presence of these redaction boxes suggests that the original document contained sensitive information that needed to be protected.



Figure 10 Foucault results from the HeNe laser

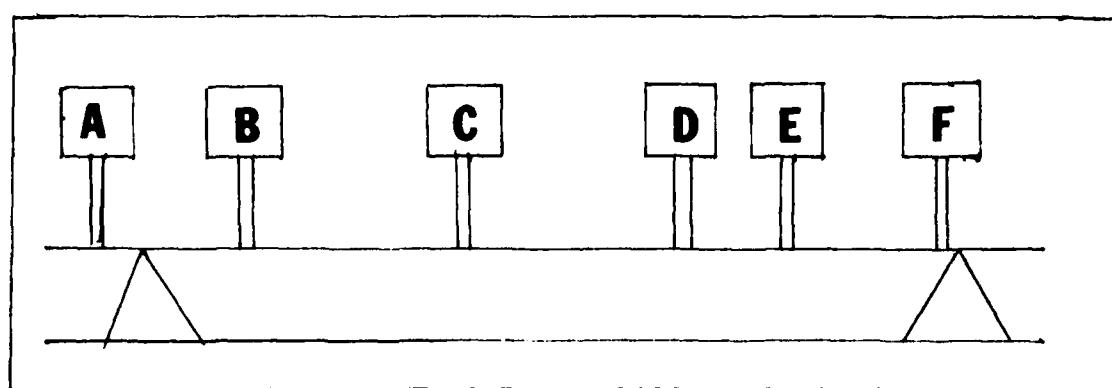


Figure 11 Experimental setup with the white light source
A is white light source
B is the pinhole
C is the lens under test
D is the knife-edge
E is the imaging lens
F is screen/camera

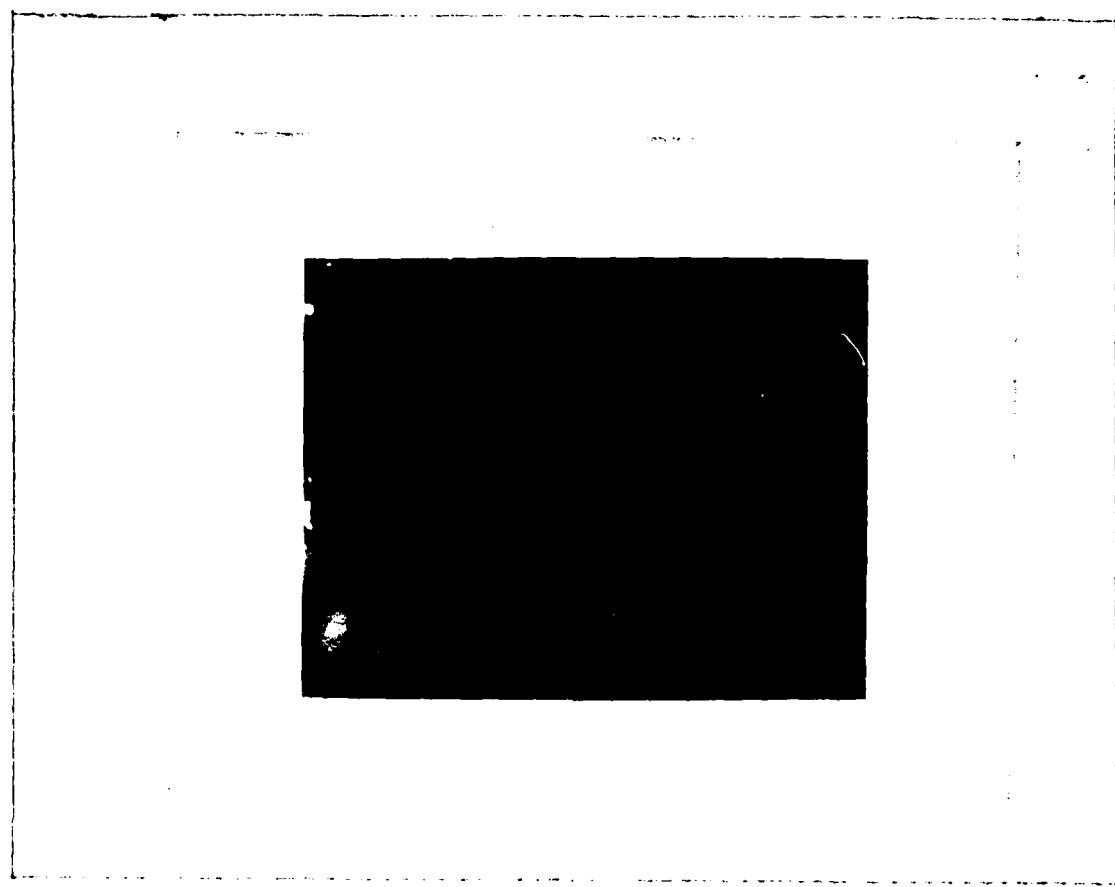


Fig. 12. Faint light source from the white light source

The doublet lens and Jelco beam expander in the setup were replaced by a white light source and a pinhole aperture (Fig. 11). The pinhole was placed in front of the objective lens to prevent interference in the intermediate image plane. The beam splitter was placed in front of the eyepiece lens to prevent interference in the final image plane.

The doublet lens and Jelco beam expander were used to obtain a large intermediate image of the light source. The beam splitter was placed in front of the eyepiece lens to prevent interference in the final image plane. The doublet lens and Jelco beam expander were used to obtain a large intermediate image of the light source. The beam splitter was placed in front of the eyepiece lens to prevent interference in the final image plane.

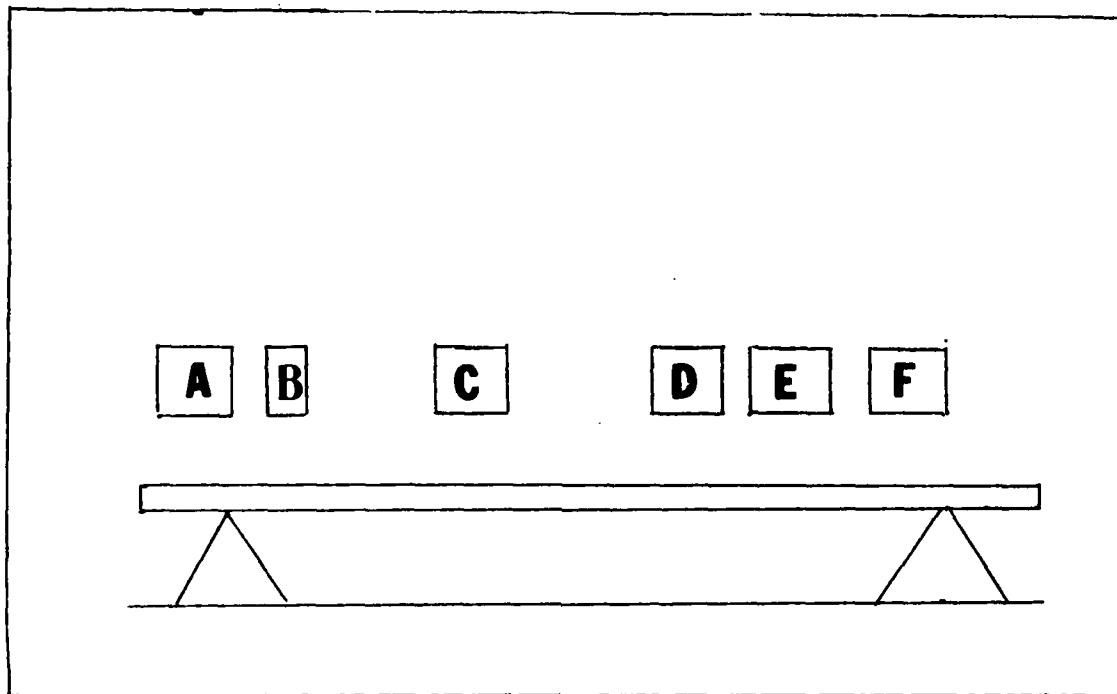


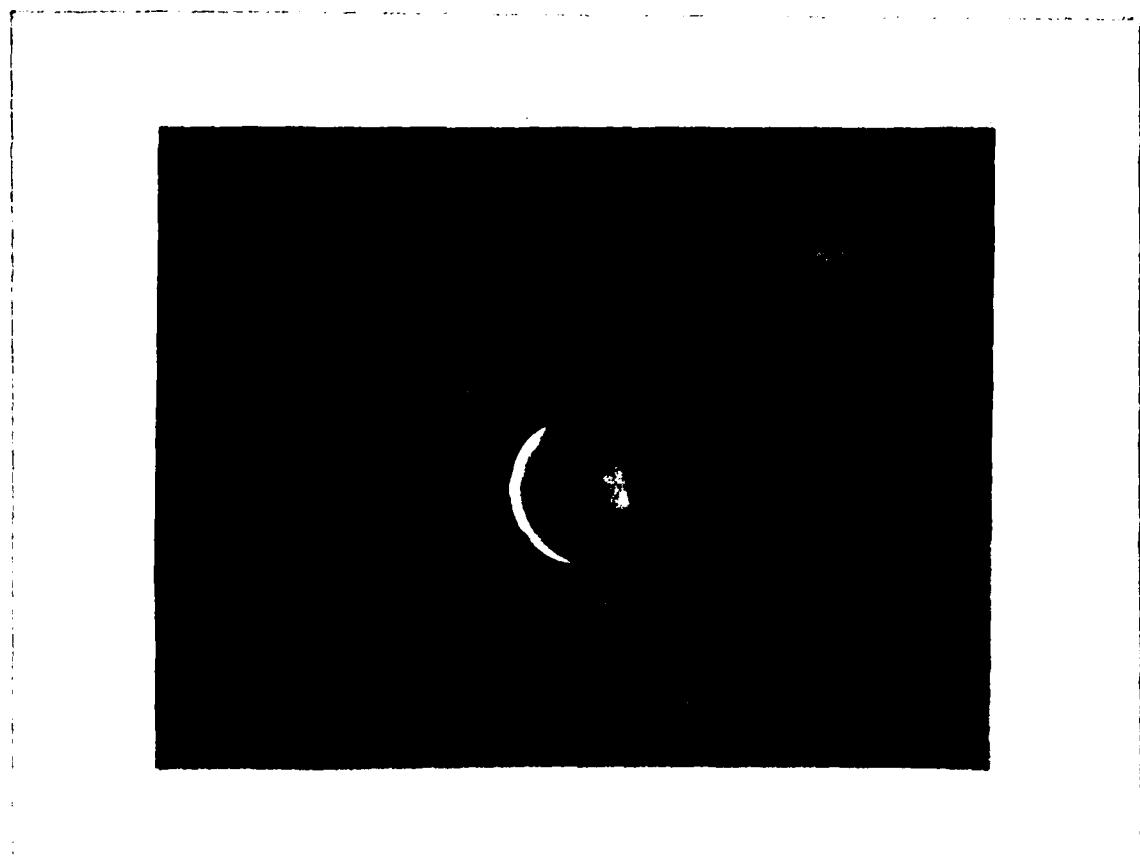
Figure 13 Experimental setup sodium source

- A is the sodium lamp
- B is the pin hole
- C is the lens under test
- D is the knife-edge
- E is the imaging lens
- F is the screen/camera

across the image at all times regardless of the amount of the beam the knife-edge cut off. The results that were obtained are presented in Figure 12.

The setup was then altered by the replacement of the white light source with a Sodium lamp as in Figure 13. The lenses' positions were adjusted for the wavelength with their order on the optic bench remaining the same. This setup gave the true Focogram that is presented as Figure 14.

A test to see how the aperture affects the image was conducted by inserting a 3 by 5 inch index card into the



40

the Foucault test. The actual explanation as to what is happening is presented in Chapter 2 pages 17-20.

Chapter 4

The Design a Ball-Hall Test

To interpret the information from the Foucault graph in an adaptive optical system, the theory that is presented in chapter two is used as the starting point. Each of the images of the knife-edge is scanned to yield the irradiance as a function of x and y . The irradiance is then divided by the background irradiance and run through a transformation to obtain the slope of the mirror surface. This slope of the mirror surface is then integrated to yield the height of the surface. Then the square of the differences of the integrated height for each direction at each point (x, y) , is minimized to obtain the local curvature of the mirror surface. The derivation of the equations to accomplish this task is as follows.

For a point in the x -direction, if the scan of the irradiance, $I(x)$, is over Δx and the center is x_0 , and the surface height, $z(x)$ is integrated over the interval Δx centered at x_0 , then $I(x_0)$ is obtained. The integrated height, up to figure 4.1, is given by

$$z(x) = \int_{x_0 - \frac{\Delta x}{2}}^{x_0 + \frac{\Delta x}{2}} \int_{y_0 - \frac{\Delta y}{2}}^{y_0 + \frac{\Delta y}{2}} \frac{I(x, y)}{I(x_0, y_0)} \sum_{i=1}^n \frac{1}{\sqrt{(x - x_i)^2 + (y - y_i)^2}}$$

Figure 4.1: A diagram showing a rectangular area with a grid of points (x_i, y_i) .

Figure 4.2: A diagram showing a rectangular area with a grid of points (x_i, y_i) .

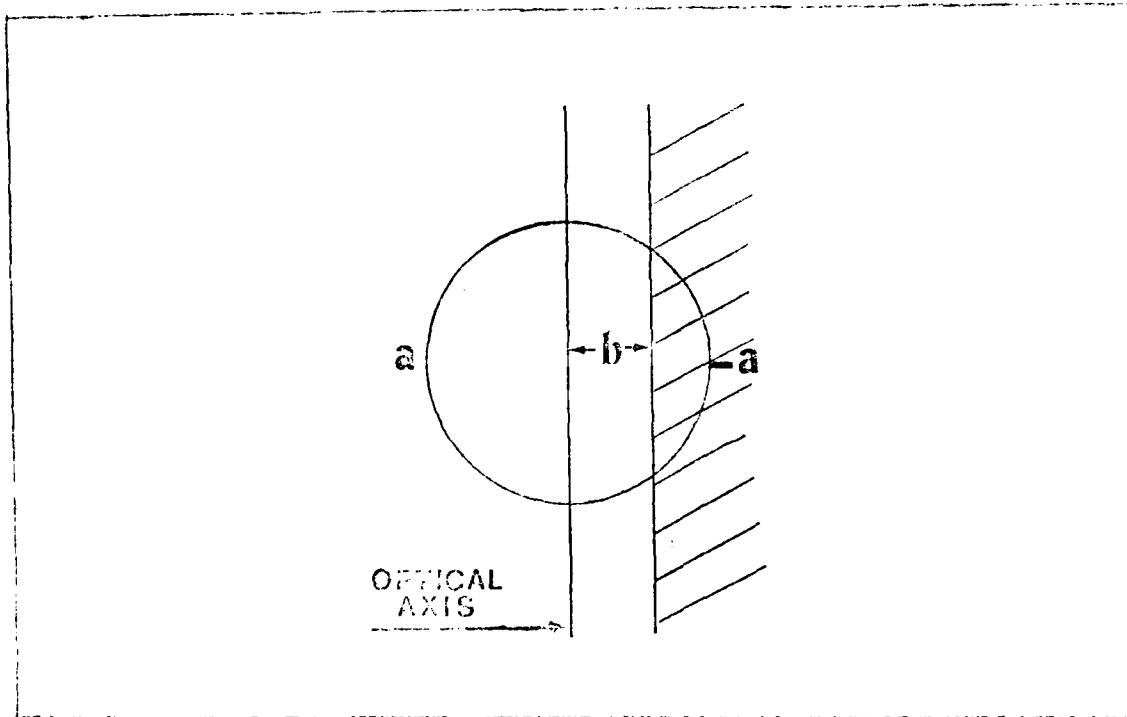


Figure 15. Method of measurement for calculating the slope of the curve.

a = distance of center

b = distance between left edge of lens and center of curvature

$$k = \frac{a^2}{b} \int_{-\pi/2}^{\theta(b/a)} d\theta \tan^2 \theta$$

$$k = \sqrt{a^2 + b^2} - a = \sqrt{b^2 + k^2}$$

$\theta = \arctan \frac{b}{a}$

$$k = \sqrt{a^2 + b^2} - a = \sqrt{b^2 + k^2}$$

$k = \sqrt{a^2 + b^2} - a = \sqrt{b^2 + k^2}$

$\theta =$

0

where $b = 2Lm(x)$

L = distance between the mirror and knife-edge

$m(x)$ = the local slope of the mirror

The irradiance, $I(x)$, is then:

$$I(x) = I_0 A$$

$$= I_0 (\pi/2 + \arcsin(b/a) + (b/a)(1 - (2Lm(x)/v)^2)^{1/2})$$

where I_0 is the background irradiance.

Repeating b with $2Lm(x)$ yields:

$$I(x)/I_0 = \pi/2 + \arcsin((2Lm(x)/a) - (2Lm(x)/a)(1 - (2Lm(x)/v)^2)^{1/2})$$

Since the y variable has been a constant all through this derivation, and implied, it can be denoted as part of the function of x :

$$I(x)/I_0 = \pi/2 + \arcsin((2Lm(x)/a) - (2Lm(x)/a)(1 - (2Lm(x)/v)^2)^{1/2})$$

Let $C = 2Lm(x)/a$, $\theta = \arcsin((2Lm(x)/a) - (2Lm(x)/a)(1 - (2Lm(x)/v)^2)^{1/2})$, $v = v$, and $\alpha = \pi/2 - \theta$. Then every thing you need to calculate the irradiance is contained in the equation:

$$I(x)/I_0 = \pi/2 + \arcsin((C - C(1 - (v^2/C)^2)^{1/2}) - C(1 - (v^2/C)^2)^{1/2} \cos(\alpha))$$

or

$I(x) = I_0$

$$m(x, y) = \partial h(x, y) / \partial x$$

$$n(x, y) = \partial h(x, y) / \partial y$$

$$h'(x, y) = \int_0^y dx' m(x', y) + f(y)$$

$$h''(x, y) = \int_0^y dy' n(x, y') + g(x)$$

The mean difference squared between the slopes for each spot, S , can only be minimized by varying the functions $f(y)$ and $g(x)$ in the equation:

$$S = \frac{f}{g} \int_0^y dx dy (m'(x, y) - h''(x, y))^2$$

$$m'(x, y) = \int_{x-y}^x m(x', y) dx'$$

$$h''(x, y) = \int_{-y}^y h'(x, y') dy'$$

$$\text{and } S = \frac{f}{g} \int_0^y dx dy \left[\int_{x-y}^x m(x', y) dx' - \int_{-y}^y h'(x, y') dy' \right]^2$$

which is zero if

$m'(x, y) = h''(x, y)$

$\therefore m(x, y) = h(x, y)$

$\therefore f(y) = g(x) = 0$

So that

$$\delta S / \delta \mu = 0 \quad \text{when evaluated at } \mu = 0$$

$$= \int_0^a dx dy (h'(x, y) - h''(x, y)) \tau(y)$$

$$\delta S / \delta \beta = 0 \quad \text{when evaluated at } \beta = 0$$

$$= \int_0^a dx dy (h'(x, y) - h''(x, y)) \mu(x)$$

$$0 = \int_0^a dy \tau(y) \int_0^a dx (h'(x, y) - h''(x, y))$$

Participate in the limits of integration:

$$0 = \int_{-\pi}^{\pi} d\varphi \tau(\varphi) \int_{-\pi}^{\pi} d\varphi' (\tau'(\varphi, \varphi') + \tau''(\varphi, \varphi'))$$

$$\tau'(\varphi, \varphi') = \frac{d}{d\varphi} \tau(\varphi') \Big|_{\varphi' = \varphi}$$

$$\tau''(\varphi, \varphi') = \int_{-\pi}^{\pi} d\varphi'' \frac{d}{d\varphi'} \tau(\varphi'') \Big|_{\varphi'' = \varphi}$$

$$\tau'(\varphi, \varphi') = \int_{-\pi}^{\pi} d\varphi'' \tau(\varphi'') \Big|_{\varphi'' = \varphi}$$

$$\tau''(\varphi, \varphi') = \int_{-\pi}^{\pi} d\varphi'' \int_{-\pi}^{\pi} d\varphi''' \tau(\varphi''') \Big|_{\varphi''' = \varphi}$$

```

9      GOSUB 2000
0  MX(I,J) = (RAD*SL(I,J))/(2*LE)
1  IF IY(I,J)=0 THEN NX(I,J)=0:GOTO 100
2  I(I,J)=IY(I,J)
3  GOSUB 2000
5  NX(I,J) = (RAD*SL(I,J))/(2*LE)
00    NEXT J
05    NEXT I
06 PRINT "Starting to integrate m(x,y)"
10 H = RAD/HOR
130 REM Intergrate mx
180 FOR A = 1 TO HOR STEP 1
190 FOR J = 1 TO VRT STEP 1
196 S=MX(A,J)
200 FOR I = 2 TO HOR STEP 2
210   FOR CON = 1 TO 2 STEP 1
220     IF CON = 1 THEN S = S + 4*MX(I,J)
230     IF CON = 2 THEN S = S + 2*MX(I+1,J)
240   NEXT CON
250   NEXT I
260   M(A,J) = H * S / 3
270   NEXT J
275 NEXT A
276 REM Intergrate n(x,y)
277 H=RAD/VRT
280 PRINT "starting to integrate n(x,y)"

```

```

14 INPUT"Input the distance between the mirror and knife-edge";LE
15 INPUT"Input the number of points in x direction"; HOR
16 INPUT"Input the number of points in y direction"; VRT
18 INPUT "Input the background intensity"; IO
19 INPUT"Input radius of the mirror";RAD
20 REM Read Intensity in x direction
21 DIM IX(HOR,VRT),IY(HOR,VRT),MX(HOR,VRT),NX(HOR+1,VRT+1),P(VRT,HOR),G(HOR+1),
   F(VRT+1),M(HOR,VRT), N(HOR,VRT), ROY(VRT),ROX(HOR),H(HOR,VRT),FF(VRT+1),SL(V
   RT,HOR),I(VRT,HOR)
25 FOR I = 1 TO HOR
30   FOR J = 1 TO VRT
35     READ IX(I,J)
40   NEXT J
45 NEXT I
50 REM Read Intensity in y direction
55 FOR I = 1 TO HOR
60   FOR J = 1 TO VRT
65     READ IY(I,J)
70   NEXT J
75 NEXT I
80 REM Make transformation from intensity to slope of mirror
81 PRINT "Starting to calculate the slopes at each point on the mirror."
85 FOR I = 1 TO HOR
86   FOR J = 1 TO VRT
87 IF IX(I,J)=0 THEN MX(I,J)=0:GOTO 91
88 I(I,J)=IX(I,J)

```

Appendix A

Computer program to find the error in the surface of a mirror

The following computer program is written in Rising Star's version of Basic for the Epson QX-10 personal computer. The program reads the irradiance from the data file. The slope of both directions is then calculated and integrated. After an arbitrary function for $g(x)$ is chosen, an iteration is then set up and run to find the overall functions of each direction. Finally, the error in height is calculated and printed out.

The data file listed is from the case where it was assumed that the mirror had an height error function of

(*)
$$f(x) = \frac{2}{\pi} \sin(\frac{\pi}{2}x/R) + C_0$$
 where $C_0 = f_0(0) = 0$ and R is the radius of the mirror. This was chosen as it had been suggested by one user

that this would be a good model for the height profile. The following is the code for the same computer for a parabolic height profile. It is the same except for the addition of a few lines to calculate the height profile across the entire width of the mirror.

checking out of mirrors in the infrared region, or examining microwave reflectors (antennae). And there are a vast number of possibilities in the visible region alone, for example, the automation of a lens/mirror grinding system, or the stabilization of a laser cavity. The basic design is the same for all cases, just a matter of proper scale.

Chapter 5

Conclusions and Recommendations

The work of this thesis has shown that all previous diffraction theory on the Foucault knife-edge test was inadequate to describe the actual results of the test. These inadequacies are: no bright ring is found around the image, no irradiance pattern is found outside of the mirror's boundary, coherent monochromatic light source does not produce a Focogram, and that it is not the field of the waves that add at the focal point but their intensity to produce a Focogram. It has been shown that from the irradiance of the Focograms, the error of the surface on the mirror can be detected and corrected. Also presented in this thesis are the design and major components comprising an adaptive optical system built around the Dallmuller telescope.

A possible future research direction is the fine-tuning of the computer code and the actual implementation into a circuit system, like that of the CDS project, already started at the end of Chapter 4. The project currently working on the LIGO detector has developed several types of detectors as they attempt to determine the source of gravitational waves by using various polarizations of the wave. One such detector is currently being developed.

The author would like to express his thanks to Dr. James E. Gossard for his guidance and support throughout this work. He would also like to thank Dr. John C. Howell for his support and encouragement.

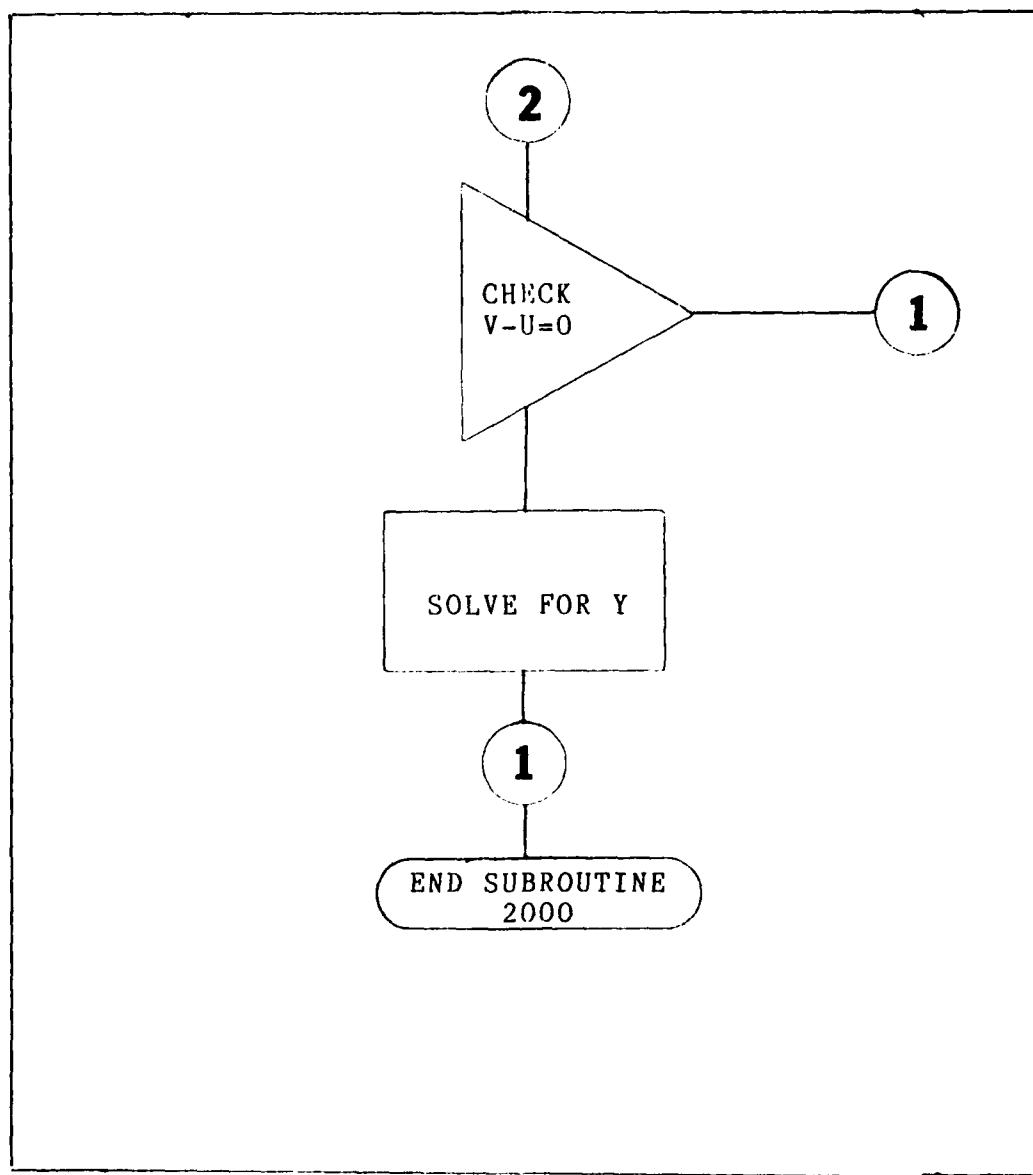


Figure 17d. The continuation of the flow chart.

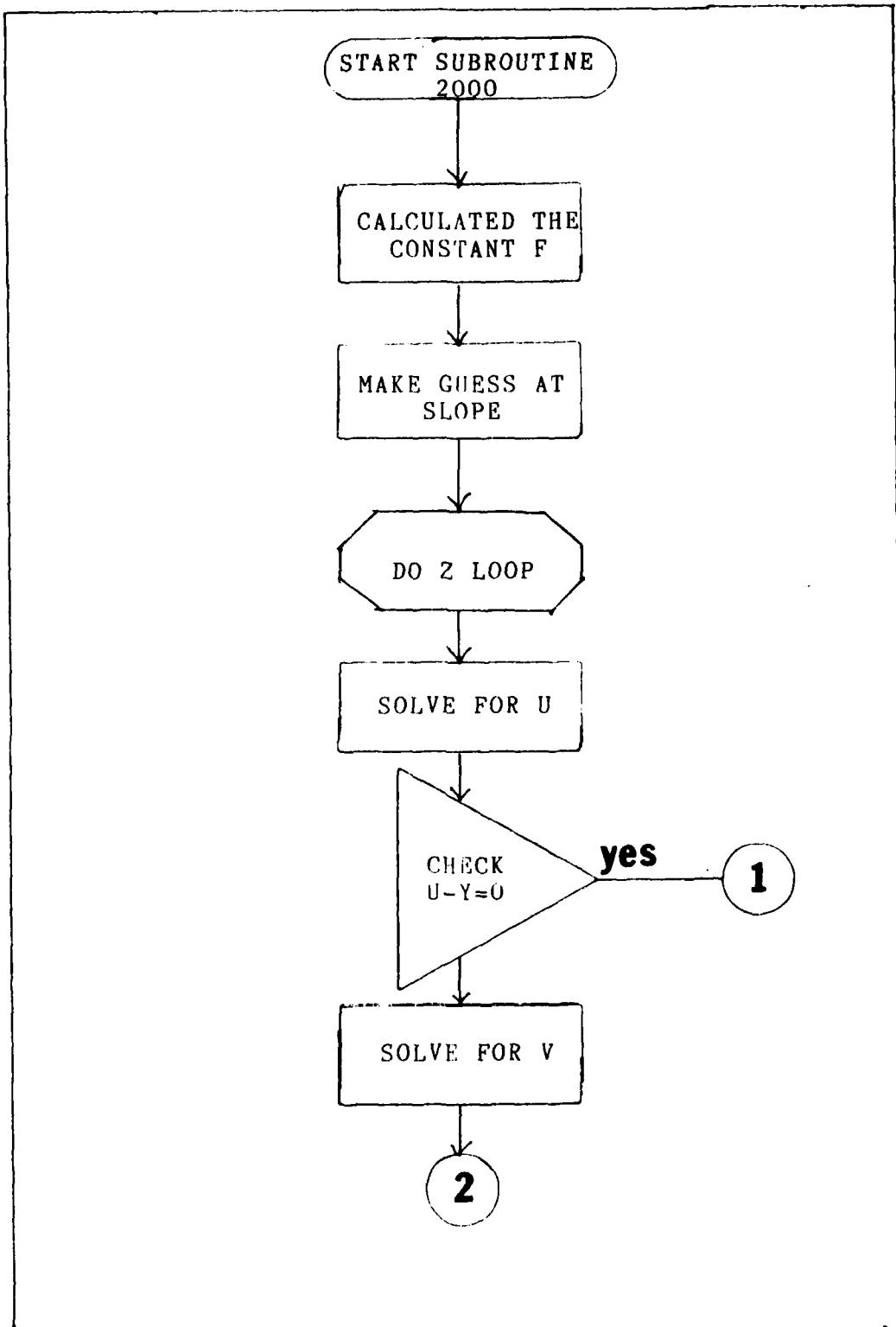


Figure 17c. The continuation of the flow chart.

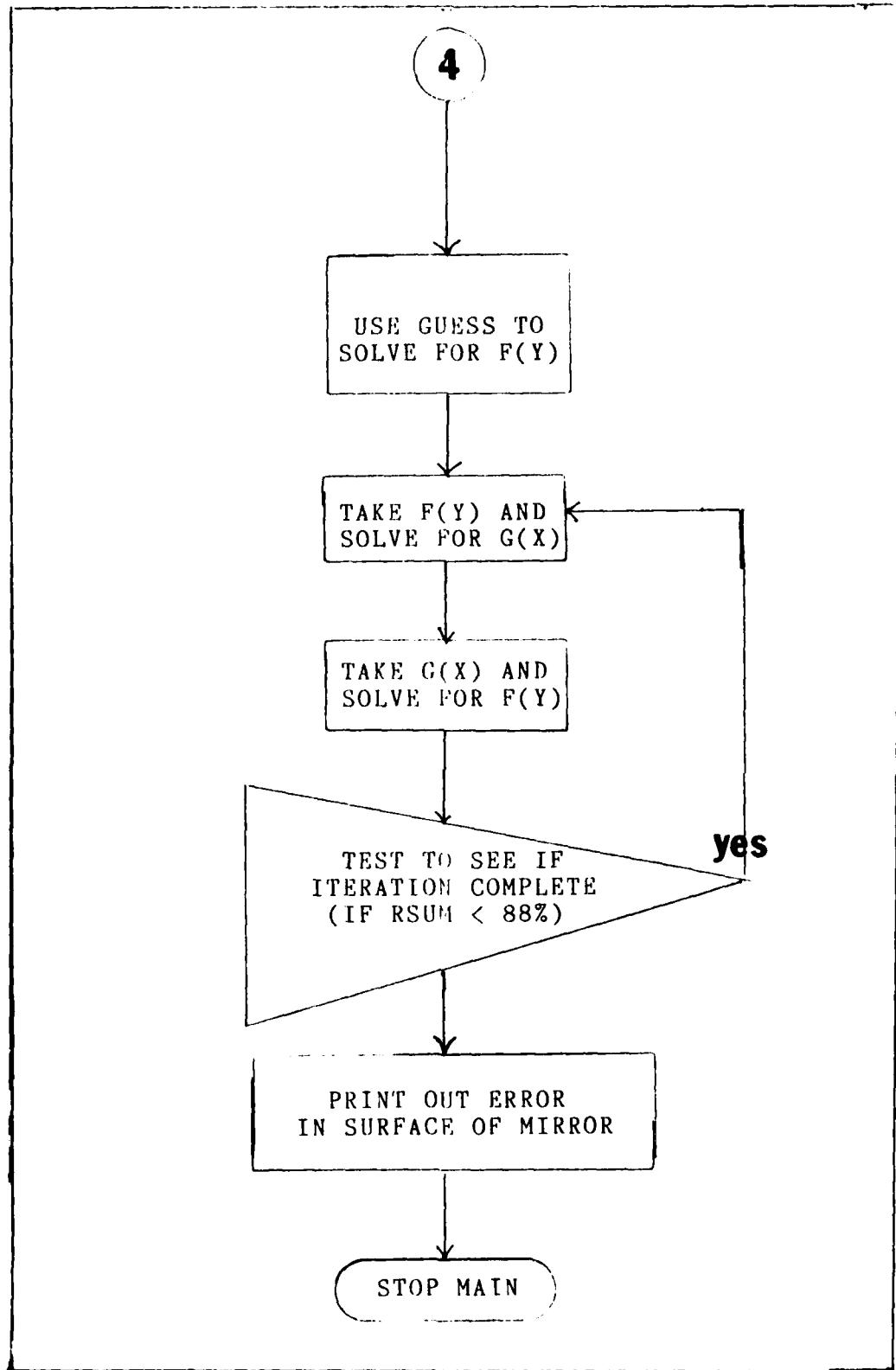


Figure 17b. The continuation of the flow chart.

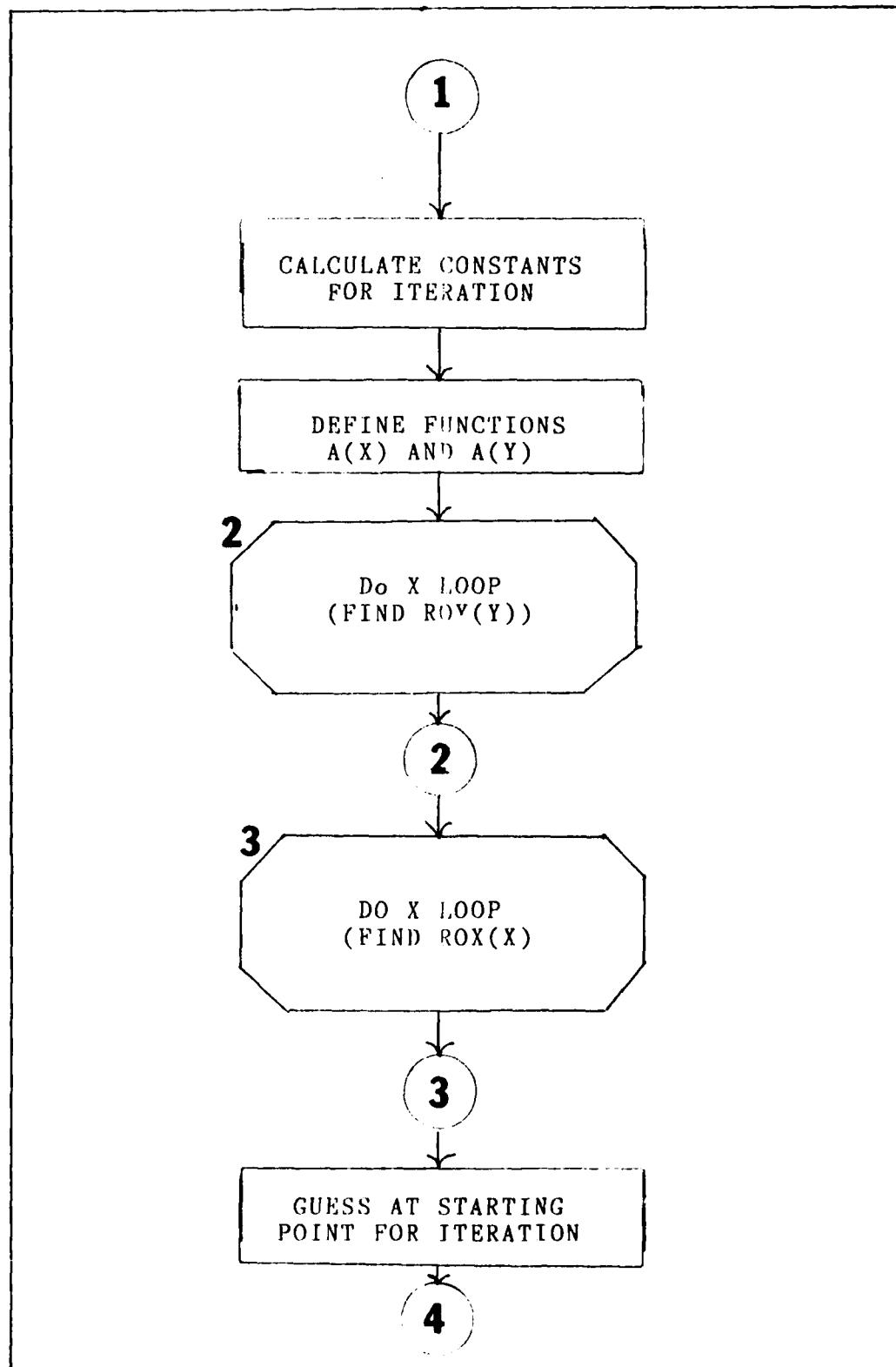


Figure 17a. The continuation of the flow chart.

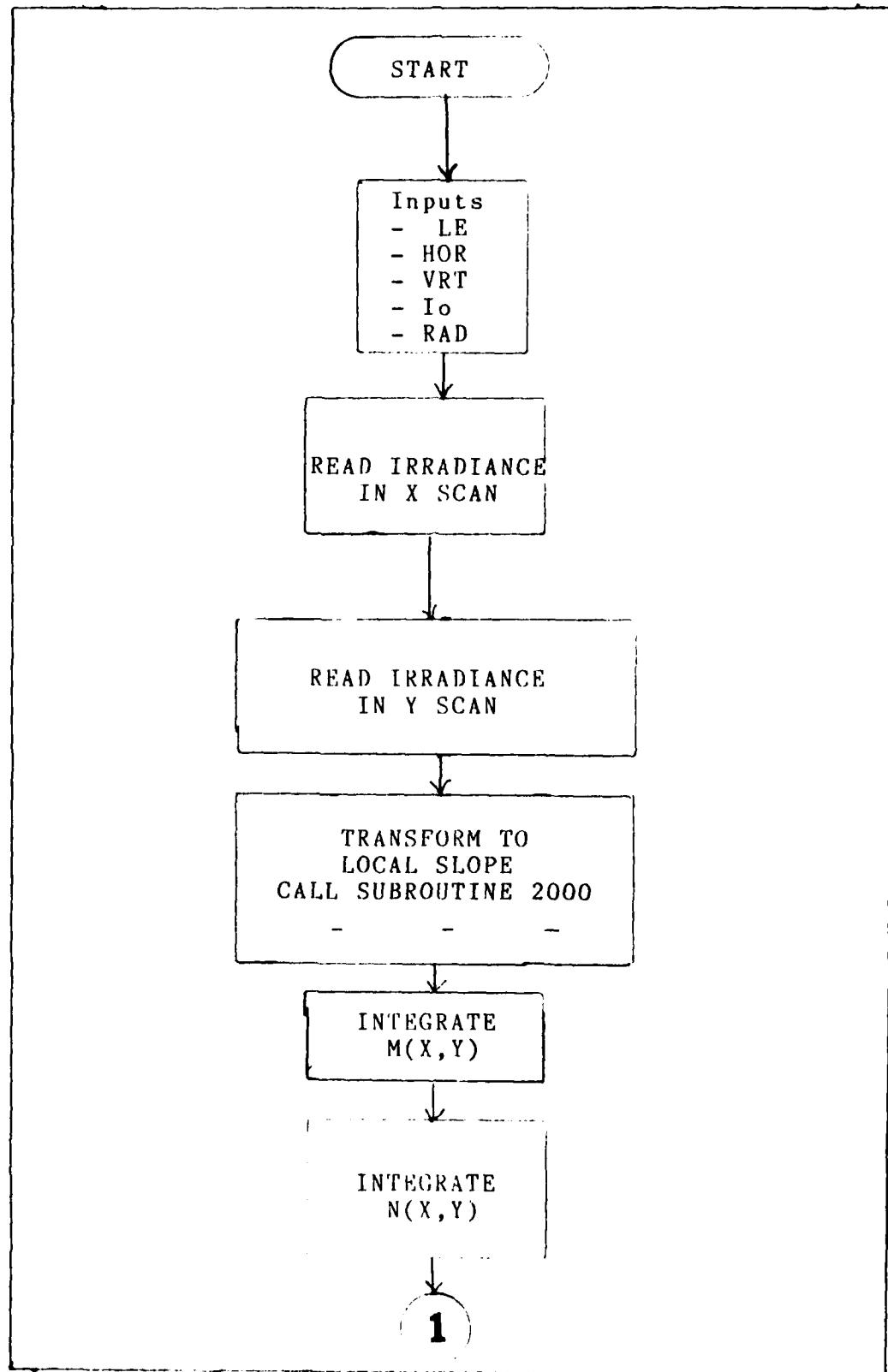


Figure 17. Flow chart of the computer program.

not using a system of this type because they are more concerned with wavefront sensing and controlling the wavefront on the out going wave, than with the surface of the mirror. They use fast acting actuators on the rear of the mirror to add imperfections to the wave front that cancel out the defects that occur during propagation to the target.

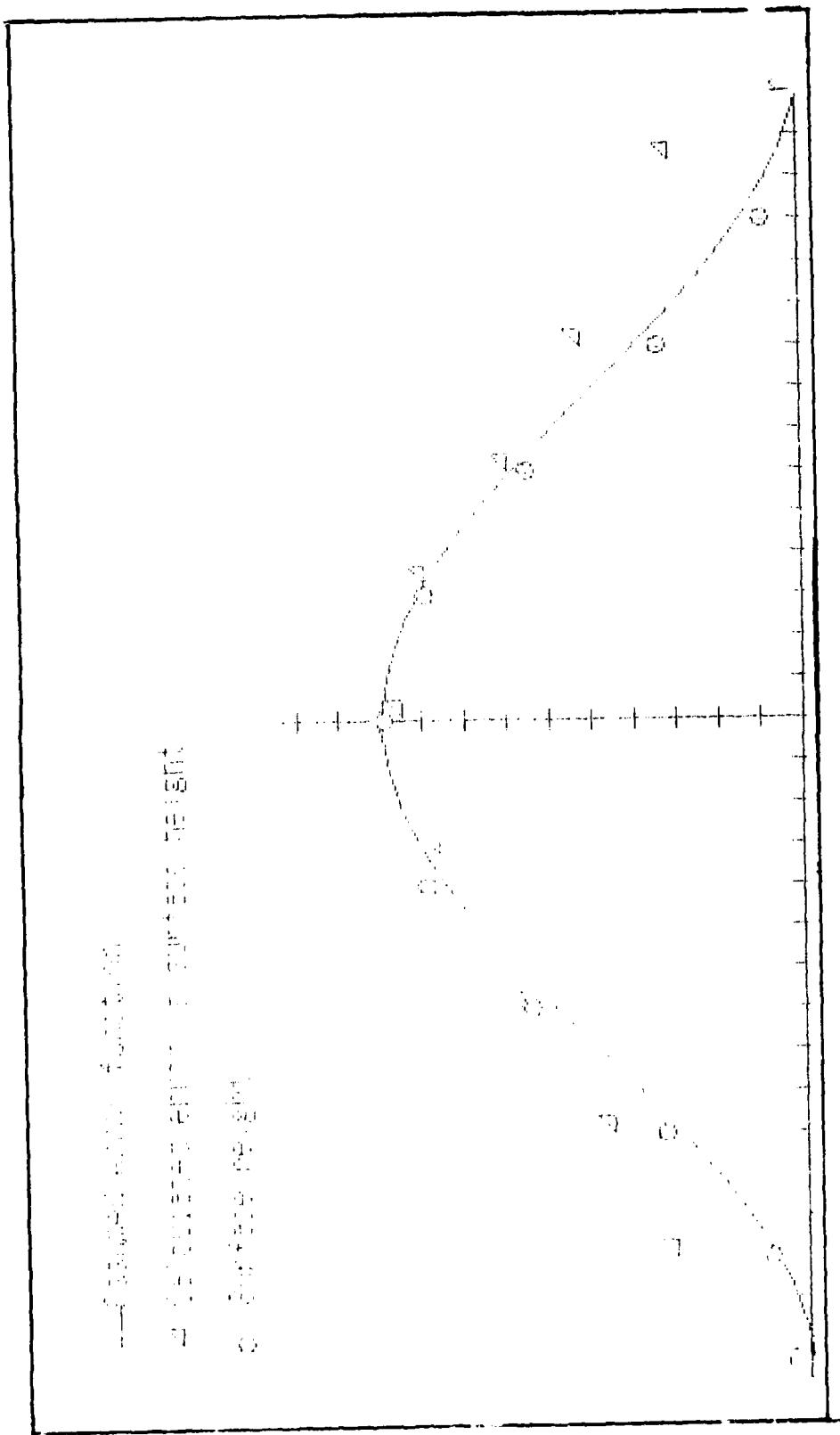


Figure 16. The graphed output across the mirror on the x-axis.

A computer program to solve for $h(x,y)$ from the above approach is included as appendix A. Appendix B is an output from a sample run. Figure 16 shows the graphed results along the x axis of a run where it was assumed the surface had an error function of $1 + \cos(\pi r/a)$, where a

was the radius and $r = (x^2 + y^2)^{1/2}$. The error function is represented as the solid curve. The height of the mirror at the points tested is shown as the circles, and the calculated error in the surface is shown as the triangles. The calculated error is close to the actual height except near the edge of the mirror. Upon closer examination, the computer code did not find (under the conditions used) that this slight variation was necessary for a good fit. This conclusion was based on the fact that the calculated error in the surface was less than 1% of the total height.

The reason for this behavior is that the computer uses the calculated error in the surface to calculate the height of the mirror. If the calculated error is small, the program will not change the height of the mirror.

The computer program can be used to calculate the height of the mirror for any given set of points. The user can enter the coordinates of the points and the radius of the mirror. The program will then calculate the height of the mirror for each point. The user can then compare the calculated height with the actual height of the mirror.

$$0 = \int_{-a(x)}^{a(x)} dy (\bar{M}(x, y) + f(y) - \bar{N}(x, y) - g(x)) \text{ for any } x$$

$$0 = 2f(y)a(y) + \int_{-a(y)}^{a(y)} dx \bar{g}(x) + \int_{-a(y)}^{a(y)} dx (\bar{M}(x, y) - \bar{N}(x, y))$$

$$0 = 2g(x)a(x) + \int_{-a(x)}^{a(x)} dy f(y) + \int_{-a(x)}^{a(x)} dy (\bar{M}(x, y) - \bar{N}(x, y))$$

$$\text{Let } p'(y) = \int_{-a(y)}^{a(y)} dx (\bar{M}(x, y) - \bar{N}(x, y)) \quad (4)$$

$$p''(x) = \int_{-a(x)}^{a(x)} dy (\bar{M}(x, y) - \bar{N}(x, y)) \quad (5)$$

$$f(y) = 1/2a(y) \left[p'(y) + \int_{-a(y)}^{a(y)} dx g(x) \right]$$

$$g(x) = 1/2a(x) \left[p''(x) + \int_{-a(x)}^{a(x)} dy f(y) \right]$$

Type III boundary conditions are given by:

$$L(\phi_{\pm}) = L(\phi_{\pm}^{\prime}) + L(\phi_{\pm}^{\prime\prime}) + L(\phi_{\pm}^{\prime\prime\prime}) + \gamma(\phi_{\pm} - \phi_{\pm}^{\prime})$$

$$L(\phi_{\pm}) = L(\phi_{\pm}^{\prime}) + L(\phi_{\pm}^{\prime\prime}) + L(\phi_{\pm}^{\prime\prime\prime}) + \gamma(\phi_{\pm} - \phi_{\pm}^{\prime}) \quad (6)$$

$$L(\phi_{\pm}) = L(\phi_{\pm}^{\prime}) + L(\phi_{\pm}^{\prime\prime}) + L(\phi_{\pm}^{\prime\prime\prime}) + \gamma(\phi_{\pm} - \phi_{\pm}^{\prime}) \quad (7)$$

$$L(\phi_{\pm}) = L(\phi_{\pm}^{\prime}) + L(\phi_{\pm}^{\prime\prime}) + L(\phi_{\pm}^{\prime\prime\prime}) + \gamma(\phi_{\pm} - \phi_{\pm}^{\prime}) \quad (8)$$

$$L(\phi_{\pm}) = L(\phi_{\pm}^{\prime}) + L(\phi_{\pm}^{\prime\prime}) + L(\phi_{\pm}^{\prime\prime\prime}) + \gamma(\phi_{\pm} - \phi_{\pm}^{\prime}) \quad (9)$$

```

290 FOR A = 1 TO VRT STEP 1
300 FOR I = 1 TO HOR STEP 1
306 S = NX(I,A)
310 FOR J = 2 TO VRT STEP 2
340     FOR CON = 1 TO 2 STEP 1
350         IF CON = 1 THEN S = S + 4 * NX(I,J)
360         IF CON = 2 THEN S = S + 2*NX(I,J+1)
370     NEXT CON
380     NEXT J
390 N(I,A) = H * S /3
395 NEXT I
396 NEXT A
397 REM Setting up constanst for the iterration phase.
400 FOR I = 1 TO HOR
410     FOR J = 1 TO VRT
420         P(I,J) = M(I,J) - N(I,J)
430     NEXT J
440     NEXT I
445 REM define functions a(y) and a(x)
450 DEF FN A(X) = SQR(RAD^2 - X^2)
460 DEF FN A(Y) = SQR(RAD^2 - Y^2)
465 REM solving the integrals for the constant functions of x and y
500 FOR Y=1 TO VRT STEP 1
505 S = P(1,Y)
510 FOR X = 1 TO (HOR/2-.5) STEP 2
530     S=S+4*P(X,Y)

```

```

540      S = S + 2*P(X+1,Y)
550      NEXT X
560      ROY(Y) = S*H/3
570      NEXT Y
580      FOR X=1 TO HOR
585      S=P(X,1)
590      FOR Y=2 TO (VRT/2-.5) STEP 2
595      S=S+4*P(X,Y)
600      S=S+2*P(X,Y+1)
610      NEXT Y
620      ROX(X)=S*H/3
630      NEXT X
640      REM Guess at what g(x) is, try 1-x^6
650      DEF FN G(X) = 1-(ABS(X))^6
660      DY = RAD/VRT
670      Y=0
680      REM Start point of iteration
690      FOR K = 1 TO VRT
700      Y=Y+DY
710      IF Y^2 => RAD^2 THEN GOTO 793
720      A=SQR(RAD^2-Y^2)
730      IN=0
740      FOR N=1 TO 11
750      X = -A + .2*A(N-1)
760      IN=IN+(.2*A)*FN G(X)
770      NEXT N

```

```
911 IF Y^2 => RAD THEN GOTO 992
920     A=SQR(RAD^2-Y^2)
930     IN=0
940     Z=INT(A/DX)
950     FOR N= 1 TO Z
960     IN=IN +DX*G(N)
970     NEXT N
980     FF(K)=(1/(2*A))*(ROY(K)+IN)
990     NEXT K
991 GOTO 997
992 FF(K)=0
993 GOTO 990
997 SUM=0
998 N=0
999 REM "Ready to check"
1000 FOR K= 1 TO HOR
1010     SUM=SUM+1
1020     DIF = FF(K)-F(K)
1030     IF DIF < 1E-03 THEN N=N+1
1040     F(K)=FF(K)
1050     NEXT K
1060     RSUM = N/SUM*100
1070     PRINT "The percent matched =",RSUM
1080     IF RSUM < 85 THEN 795
1090 REM if matched then print out f(y) and g(x) and find h(x,y)
1091 FOR X = 1TO HOR
```

```
780 F(K)=(1/(2*A))*(ROY(K)+IN)
785 W=W
790 NEXT K
791 GOTO 795
793 F(K)= 0
794 GOTO 790
795 DX = RAD/HOR
796 X=0
799 REM "ready to start 800 loop"
800 FOR K = 1 TO HOR
810 X=X+DX
811 IF X^2 => RAD^2 THEN GOTO 893
810 A=SQR(RAD^2 - X^2)
830 IN=0
840 Z=INT(A/DY)
850 FOR N = 1 TO Z
860 IN =IN + DY*F(Z)
870 NEXT N
880 G(K) = 1/((2*A))*(ROX(K)+IN)
890 NEXT K
891 GOTO 899
893 G(K) = 0
894 GOTO 890
899 Y=0
900 FOR K=1 TO VRT
900 Y=Y+DY
```

```

1092      PRINT F(X),G(X)
1093      NEXT X
1098      I=-1
1100      FOR X= 1 TO HOR
1105      J=-1
1110      FOR Y= 1 TO VRT
1115      IF (J^2+I^2) > RAD THEN H(X,Y)=0:GOTO 1130
1120      H(X,Y)=.5*(M(X,Y)+N(X,Y)+F(Y)+G(X))
1130      LPRINT X,Y,"Error is" H(X,Y)
1131      J=J+(2*(RAD/VRT))
1140      NEXT Y
1145      I=I+(2*(RAD/HOR))
1150      NEXT X
1160      GOTO 9999
1990 REM This is the subroutine to numericly solve for the slope of the mirror.
2000 F = I(I,J)/IO - PI/8
2100 Y=1E-03
2120 FOR Z = 1 TO 20
2200 U = SIN(F + SQR(ABS(1-Y*Y)))
2205 IF (U-Y) = 0 THEN GOTO 2300
2210 Y = 1/(U-Y)
2211 V=SIN(F+U*SQR(1-U*U))
2212 IF (V-U) = 0 THEN V=Y; GOTO 2300
2213 V=1/(V-U)
2214 IF (V-Y) = 0 THEN GOTO 2300
2215 Y=U+1/(V-Y)

```

```

220 NEXT Z

2300 SL(I,J)=Y

2310 RETURN

3000 DATA 0,0,0,0,0,0,0,0,0,0,0,0,0,.914966,1.59555,1.84658,1.59556,.914966,0
      ,0

3010 DATA 0,0,0,1.01765,2.00827,2.72604,2.98783,2.72604,2.00827,1.01765,0,0,0,.
      45748,1.33885,2.1743,2.77137,2.98783,2.77137,2.17431,1.33885,.45748,0,0,.
      39889,.90868,1.3856,1.72,1.847,1.72423,1.3857,.90868,.39889,0,0,0,0,0,0,0
      ,0,0,0,0,-.3989,-.9087,-1.3857

3020 DATA -1.72422,-1.846,-1.724,-1.3857,-.90868,-.39889,0,0,-.45748,-1.3388,-2
      .1743,-2.7714,-2.9878,-2.771,-2.174,-1.3388,-.45748,0,0,0,-1.0176,-2.0083,
      -2.726,-2.9878,-2.726,-2.0083,-1.01765,0,0,0,0,-.91497,-1.5955,-1.84658,
      -1.5955,-.91497,0,0

3030 DATA 0,0,0,0,0,0,0,0,0,0,0,0,0,0,0,0,0,0,0,0,0,0,.457483,.398889,0
      ,-.398889,-.457483,0,0,0,0,0,1.01765,1.33884,.90868,0,-.90868,-1.33884,-1.
      01765,0,0,0,.914966,2.00826,2.17431,1.38568,0,-1.38568,-2.17431,-2.00826,-
      .914966,0,0,1.59555,2.72604

3040 DATA 2.77137,1.72422,0,-1.72422,-2.77137,-2.72604,-1.59555,0,0,1.84658,2.9
      8783,2.98783,1.84658,0,-1.84658,-2.98783,-2.98783,-1.84658,0,0,1.59555,2.7
      2604,2.77137,1.72422,0,-1.72422,-2.77137,-2.72064,-1.59555,0,0,.914966,2.0
      0826,2.17431,1.38568,0,-1.38568

3050 DATA -2.17431,-2.00826,-.914966,0,0,0,1.017654,1.338884,.90868,0,-.90868,-
      1.33884,-1.01765,0,0,0,0,.457483,.398889,0,-.398889,-.457483,0,0,0,0,0,0
      ,0,0,0,0,0,0,0

9999 END

```

APPENDIX B A SAMPLE OUTPUT FROM PROGRAM

	1	Error is 0
1	2	Error is 0
1	3	Error is 0
1	4	Error is 0
1	5	Error is 0
1	6	Error is 0
1	7	Error is 0
1	8	Error is 0
1	9	Error is 0
1	10	Error is 0
1	11	Error is 0
2	1	Error is 0
2	2	Error is 0
	3	Error is 0
2	4	Error is 1.1384639849
2	5	Error is 1.1372537323
2	6	Error is 1.1388195447
2	7	Error is 1.137878966
2	8	Error is 1.1374952983
2	9	Error is 1.1380261465
2	10	Error is 0
2	11	Error is 0
3	1	Error is 0
3	2	Error is 0
3	3	Error is 1.1388024637
3	4	Error is 1.1396093156

3 5 Error is 1.1383997087
3 6 Error is 1.1399658315
3 7 Error is 1.1390248203
3 8 Error is 1.1386402233
3 9 Error is 1.1391705686
3 10 Error is 1.1364436818
3 11 Error is 0
4 1 Error is 0
4 2 Error is 1.1399686284
4 3 Error is 1.1405416143
4 4 Error is 1.1413486153
4 5 Error is 1.140138311
4 6 Error is 1.1417042413
4 7 Error is 1.1407631514
4 8 Error is 1.140378128
4 9 Error is 1.1409084712
4 10 Error is 1.1381820315
4 11 Error is 1.1151828256
5 1 Error is 0
5 2 Error is 1.1407981127
5 3 Error is 1.1413812201
5 4 Error is 1.1421781141
5 5 Error is 1.140966361
5 6 Error is 1.1425316291
5 7 Error is 1.1415906985
5 8 Error is 1.141205673
5 9 Error is 1.1417364446

5	10	Error is 1.1390109322
5	11	Error is 1.1160117263
6	1	Error is 0
6	2	Error is 1.1442853743
6	3	Error is 1.1448593163
6	4	Error is 1.1456654273
6	5	Error is 1.1444530782
6	6	Error is 1.146018696
6	7	Error is 1.1450772066
6	8	Error is 1.1446920279
6	9	Error is 1.1452228761
6	10	Error is 1.1424977963
6	11	Error is 1.1194985905
7	1	Error is 0
7	2	Error is 1.1402571981
7	3	Error is 1.1408306952
7	4	Error is 1.1416367275
7	5	Error is 1.1404245377
7	6	Error is 1.1419895967
7	7	Error is 1.1410488752
7	8	Error is 1.1406642864
7	9	Error is 1.1411953954
7	10	Error is 1.1384700051
7	11	Error is 1.1154707992
8	1	Error is 0
8	2	Error is 1.1388769069

8 3 Error is 1.1394494872
8 4 Error is 1.1402550931
8 5 Error is 1.1390429011
8 6 Error is 1.140607807
8 7 Error is 1.1396676754
8 8 Error is 1.1392846058
8 9 Error is 1.1398163441
8 10 Error is 1.13709031
8 11 Error is 1.1140911041
9 1 Error is 0
9 2 Error is 1.140134488
9 3 Error is 1.1407066998
9 4 Error is 1.1415123037
9 5 Error is 1.1403005402
9 6 Error is 1.1418655226
9 7 Error is 1.1409256518
9 8 Error is 1.1405432114
9 9 Error is 1.1410748048
9 10 Error is 1.138347918
9 11 Error is 1.1153487121
10 1 Error is 0
10 2 Error is 0
10 3 Error is 1.1330263727
10 4 Error is 1.1338326099
10 5 Error is 1.1326217736
10 6 Error is 1.1341871886

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VITA

Captain Randy L. Wingler was born on 3 April 1958 in Shell Lake, Wisconsin. He graduated from Spooner High School, Spooner, Wisconsin, in 1976 and attended the University of Wisconsin-Superior from which he received the degree of Bachelor of Science in Physics in May 1980. Upon graduation he received a commission in the USAF through the ROTC program. He was then called to active duty Aug 80 and reported to Hill AFB, Utah. There he served two and half years as the manager of the Photoreconnaissance Engineering Test Facility until entering the School of Engineering, Air Force Institute of Technology, WPAFB, Ohio, in June 1983.

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Abstract

This is a study to design a self correcting primary mirror system for a space telescope. The design is centered around a Dall-Null tester (a Foucault knife-edge tester with compensating lens). An indepth study of the theory of the Foucault test from Foucault's original publications to current work is presented. Also short comings of the diffraction approach are shown. The findings of an simple experiment showed the way to the correct explanation as to the workings of the test. Based on this new explanation, a computer program to find the error in the surface of the mirror from the irradiance pattern provided by the Dall-Null tester was developed. The computer program with a sample run is included in the appendixes A and B.

The basic design of an adaptive optic system for a space-borne application is also presented in the paper. This design has the desired quality of being able to correct the mirror while the telescope is in use. The equations being independent of wavelength allows for the design to be applied to systems working outside of the visible spectrum as well as the systems working in the visible.

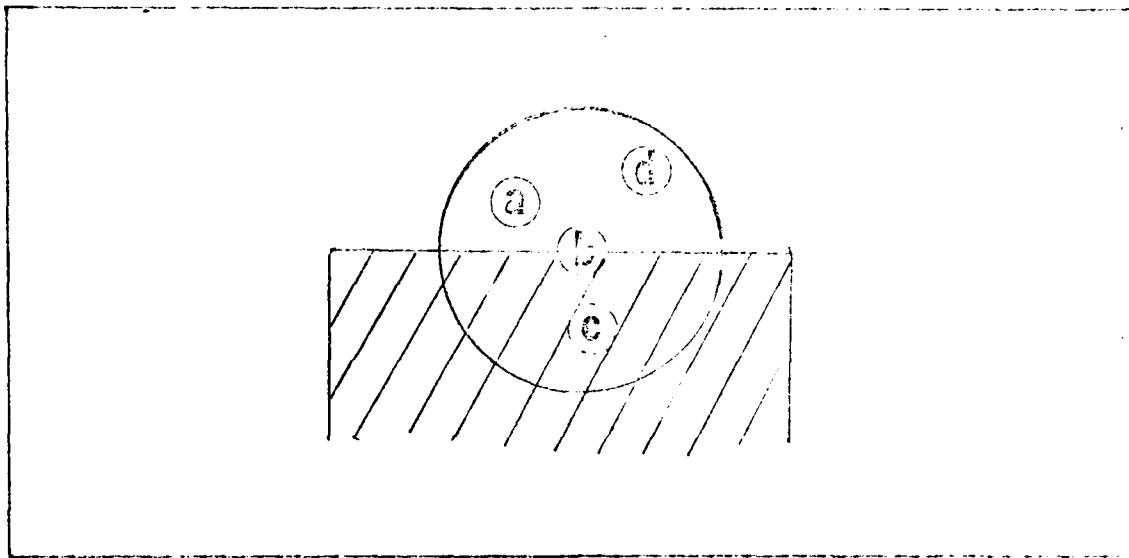


Figure 6. The relative position of the frisbee pattern at the knife-edge.

image, cases a and d, then all of the image will pass the knife-edge. But if the slope lowered the image, case c, the knife-edge would cut it off completely. Thus the boundaries of the resulting image would be determined by how much of each segment's image was passed by the knife-edge. Figure 7 shows the result for the current situation.

Thus the correct theory of the Foucault test is complete, and likewise, so is the theory of the Bell-Lulli test. And in chapter four this theory will be put to work to find the error in the height of the mirror surface.

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